

Six- and Eight-Vertex models on their combinatorial line

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Spin chains

A standard model for magnetism in one dimension:

$$H = -\frac{1}{2} \sum_{i=1}^L (J_x \sigma_i^x \sigma_{i+1}^x + J_y \sigma_i^y \sigma_{i+1}^y + J_z \sigma_i^z \sigma_{i+1}^z)$$

where the $\sigma_i^{x,y,z}$ are Pauli matrices acting on the i^{th} factor of $\mathcal{H} = (\mathbb{C}^2)^{\otimes L}$. (we assume periodic boundary conditions, i.e., $L+1 \equiv 1$)

- the XXX chain: $J_x = J_y = J_z$ [Heisenberg, 1928; Bethe, 1931]
- the XXZ chain: $J_x = J_y \neq J_z$ [Yang and Yang, 1966]
- the XYZ chain: $J_x \neq J_y \neq J_z$ [Baxter, 1973]

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Quantum integrability

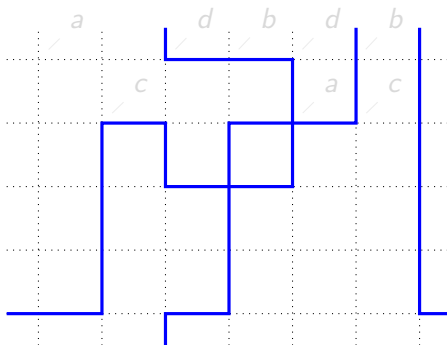
The XXX/XXZ/XYZ Hamiltonian commutes with a one-parameter family of operators on \mathcal{H} called **transfer matrices**:

$$[T(u), T(v)] = 0 \quad [H, T(u)] = 0$$

For the XXX/XXZ case, $u \in \mathbb{C}$, whereas in the XYZ case, u lives on an elliptic curve.

Six- and Eight-vertex models

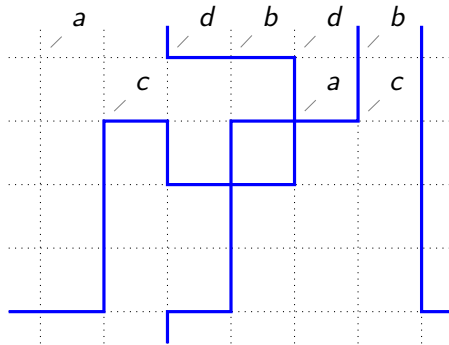
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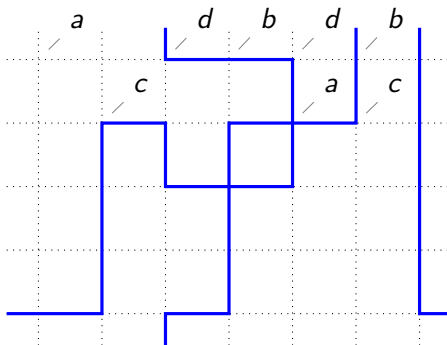
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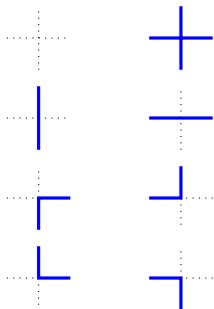


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Transfer matrix

The operator $T(u)$ has the meaning of row-to-row transfer matrix on the two-dimensional square lattice: $\mathbb{C}^2 = \langle | = +, | = - \rangle$

$$\langle \alpha' | T(u) | \alpha \rangle = \sum_{\beta_1, \dots, \beta_L = \text{---}, \dots} \frac{\alpha'_1 \quad \alpha'_2 \quad \dots \quad \alpha'_L}{\beta_1 \quad \beta_2 \quad \beta_3 \quad \dots \quad \beta_L \quad \beta_1} \frac{\alpha_1 \quad \alpha_2 \quad \dots \quad \alpha_L}{\beta_1 \quad \beta_2 \quad \beta_3 \quad \dots \quad \beta_L \quad \beta_1}$$



$$a(u) = \vartheta_4(2\eta)\vartheta_4(u)\vartheta_1(u + 2\eta)$$

$$b(u) = \vartheta_4(2\eta)\vartheta_1(u)\vartheta_4(u + 2\eta)$$

$$c(u) = \vartheta_1(2\eta)\vartheta_4(u)\vartheta_4(u + 2\eta)$$

$$d(u) = \vartheta_1(2\eta)\vartheta_1(u)\vartheta_1(u + 2\eta)$$

Ground state

We are particularly interested in the **ground state** of the spin chain, that is the eigenvector of H corresponding to its lowest eigenvalue. In a certain range of u it coincides with the largest eigenvalue of $T(u)$.

If L is odd the ground state is two-fold degenerate, but we make it unique by imposing $\prod_{i=1}^L \sigma_i^x = 1$.

In general, using the **Bethe Ansatz** (or other related methods), one can obtain explicit formulae for the ground state only in the thermodynamic limit where $L \rightarrow \infty$.

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Conjecture (Baxter, '72,'73,'89; Stroganov, '99)

Suppose L is odd, and

$$J_x J_y + J_x J_z + J_y J_z = 0$$

Then the ground state eigenvalue is

$$E_0 = -\frac{L}{2}(J_x + J_y + J_z)$$

Combinatorial line

Parameterize (up to normalization)

$$J_x = 1 + \zeta \quad J_y = 1 - \zeta \quad J_z = \Delta$$

then $J_x J_y + J_x J_z + J_y J_z = 0$ equivalent to $\Delta = \frac{\zeta^2 - 1}{2}$.

Further,

$$\zeta = \left(\frac{\vartheta_1(2\eta, p)}{\vartheta_4(2\eta, p)} \right)^2 \quad \Delta = \frac{\vartheta_4^2(0)}{\vartheta_2(0, p)\vartheta_3(0, p)} \frac{\vartheta_2(2\eta, p)\vartheta_3(2\eta, p)}{\vartheta_4^2(2\eta, p)}$$

then $J_x J_y + J_x J_z + J_y J_z = 0$ equivalent to $\eta = \pi/3$.

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then $J_x J_y + J_x J_z + J_y J_z = 0$ equivalent to $\eta = \pi/3$.

XXZ/6-vertex case

The XXZ/6-vertex case corresponds to $\zeta = 0$, or $p = 0$. The condition on the J 's means

$$H = -\frac{1}{2} \sum_{i=1}^L (\sigma_i^x \sigma_{i+1}^x + \sigma_i^y \sigma_{i+1}^y + \Delta \sigma_i^z \sigma_{i+1}^z) \quad \Delta = -\frac{1}{2}$$

In this case the conjecture above is a theorem [Yang, Fendley, '04; Veneziano and Wosiek, '06; Razumov, Stroganov and Z-J, '07].

Set $L = 2n + 1$. Call $\Psi \in \mathcal{H}$ the ground state eigenvector. Normalize it so that its entries are *polynomials* in ζ with integer coefficients and no common factors.

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Theorem (Di Francesco, Z-J and Zuber, '06; Razumov, Stroganov and Z-J, '07)

Set $\zeta = 0$ (XXZ case). We consider only entries which have n or $n + 1$ $-$'s depending on the parity of n (the other ones are zero).

- All entries are positive integers.

- $\Psi_{\underbrace{+\dots+}_n \pm \underbrace{-\dots-}_n} = 1.$

- $\Psi_{\pm \underbrace{+\dots+}_n} = A_n = \prod_{i=0}^{n-1} \frac{(3i+1)!}{(n+i)!}.$

- $\sum_{\alpha_1, \dots, \alpha_L} \Psi_{\alpha_1, \dots, \alpha_L} = \left(\frac{3}{2}\right)^n \frac{2 \times 5 \times \dots \times (3n-1)}{1 \times 3 \times \dots \times (2n-1)} A_n.$

where $\pm = (-1)^n$.

A_n is the number of **Alternating Sign Matrices** of size n .

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Conjecture (Razumov and Stroganov, '09)

- All entries of Ψ have **positive** coefficients as polynomials in ζ .
- $\Psi_{+\dots+} = \zeta^{\frac{n(n+1)}{2}} + \dots + A_V(n+1)^2 \zeta^{\lfloor \frac{n}{2} \rfloor}$.

Conjecture (Bazhanov and Mangazeev, '09)

The norm of Ψ is given by

$$|\Psi|^2 = \sum_{\alpha} \Psi_{\alpha_1, \dots, \alpha_L}^2 = (4/3)^n \zeta^{n(n+1)} s_n(\zeta^{-2}) s_{-n-1}(\zeta^{-2})$$

where the $s_n(z)$ are determined by $s_0 = s_1 = 1$ and the recurrence

$$2z(z-1)(9z-1)^2(s_n s_n'' - s_n'^2) + 2(3z-1)^2(9z-1)s_n s_n' + 8(2n+1)^2 s_{n+1} s_{n-1} - [4(3n+1)(3n+2) + (9z-1)n(5n+3)] s_n^2 = 0$$

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Inhomogeneous transfer matrix

Define a more general transfer matrix:

$$T(u; x_1, \dots, x_L) = \begin{array}{c} | \quad | \quad \dots \quad | \\ \hline u-x_1 \quad u-x_2 \quad \dots \quad u-x_L \\ \hline | \quad | \quad \dots \quad | \end{array}$$

where the weight at vertex i is a function of $u - x_i$.

Conjecture (Razumov and Stroganov, '09)

If $\eta = \pi/3$ and L is odd, then the inhomogeneous transfer matrix $T(u; x_1, \dots, x_L)$ has the eigenvalue

$$\prod_{i=1}^L (a(u - x_i) + b(u - x_i))$$

NB: in the XXZ case, this is a theorem [R, S, and Z-J, '07]; see also [DF and Z-J, '05].

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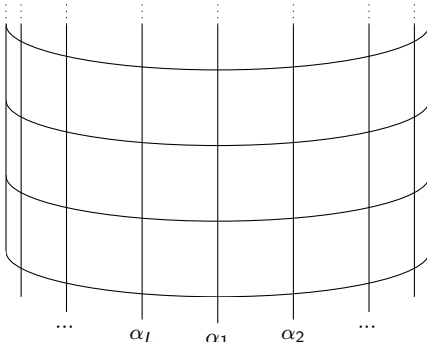
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Eigenvector

Consider the eigenvector $\Psi_L(x_1, \dots, x_L)$ of $T(u; x_1, \dots, x_L)$ with the eigenvalue $\prod_{i=1}^L (a(u - x_i) + b(u - x_i))$.

$$\Psi_{\alpha_1, \dots, \alpha_L} =$$


The diagram shows a cylinder with several vertical lines representing poles. The poles are labeled with parameters $\alpha_1, \alpha_2, \dots, \alpha_L$. The surface of the cylinder is divided into horizontal bands by curved lines.

Relation to quantum Knizhnik–Zamolodchikov

- $\Psi(x_1, \dots, x_L)$ satisfies the **exchange relation**

$$\check{R}_i(x_{i+1} - x_i)\Psi(x_1, \dots, x_L) = \Psi(x_1, \dots, x_{i+1}, x_i, \dots, x_L)$$

$$\text{where } \check{R}_i(x) = \frac{1}{a(x)+b(x)} \begin{pmatrix} a(x) & 0 & 0 & d(x) \\ 0 & c(x) & b(x) & 0 \\ 0 & b(x) & c(x) & 0 \\ d(x) & 0 & 0 & a(x) \end{pmatrix} \quad ; \quad i, i+1$$

- Cyclic invariance

$$\Psi_{\alpha_1, \dots, \alpha_L}(x_1, \dots, x_L) = \Psi_{\alpha_2, \dots, \alpha_L, \alpha_1}(x_2, \dots, x_L, x_1)$$

These can be considered as a special case of the level 1 quantum Knizhnik–Zamolodchikov(–Bernard) equation at a cubic root of unity.

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Six-vertex case

We assume L odd, $\Delta = -\frac{1}{2} = \frac{1}{2}(q + q^{-1})$, $q = e^{2\pi i/3}$. The normalization of Ψ_L can be chosen so that its entries are **polynomials** (of degree $n - 1$) in each variable.

Theorem (R, S and Z-J, '07)

If $\Phi(x) = (\Phi_+(x), \Phi_-(x))$ is the type I vertex operator associated to level 1 highest weight modules of $U_q(\widehat{sl}(2))$, then $\Psi(x_1, \dots, x_L)$ coincides (up to normalization) with $\langle 0 | \Phi(x_1) \cdots \Phi(x_L) | 1 \rangle$.
 Explicitly, if $\{i : \alpha_i = +\} = \{a_1 < \cdots < a_n\}$, and $z_k = e^{2ix_k}$,

$$\Psi_{a_1, \dots, a_n}(x_1, \dots, x_L) = \prod_{1 \leq i < j \leq L} (q z_i - q^{-1} z_j)$$

$$\oint \prod_{\ell=1}^n \frac{dw_\ell}{2\pi i} \prod_{\ell=1}^n x_{a_\ell} \frac{\prod_{\ell=1}^n w_\ell \prod_{1 \leq \ell < m \leq n} [(w_m - w_\ell)(q w_\ell - q^{-1} w_m)]}{\prod_{\ell=1}^n \left[\prod_{1 \leq i \leq a_\ell} (w_\ell - z_i) \prod_{a_\ell \leq i \leq L} (q w_\ell - q^{-1} z_i) \right]}$$

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Six-vertex cont'd: linear sum rule

For example, one can compute

$$\sum_{\alpha_1, \dots, \alpha_L} \Psi_{\alpha_1, \dots, \alpha_L}(x_1, \dots, x_L) = s_{(n-1, n-1, \dots, 1, 1)}(z_1, \dots, z_L) \quad z_k = e^{2ix_k}$$

This is closely related to the enumeration of ASMs. In fact, in the six-vertex case, there is an analogue problem in even size (twisted XXZ chain). There again,

$$\sum_{\alpha_1, \dots, \alpha_{2n}} \Psi_{\alpha_1, \dots, \alpha_{2n}}(x_1, \dots, x_{2n}) = s_{(n-1, n-1, \dots, 1, 1)}(z_1, \dots, z_{2n})$$

and it coincides with the partition function of the six-vertex model with Domain-Wall Boundary Conditions, whose configurations are in bijection with Alternating Sign Matrices.

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Eight-vertex case

Go back to the general 8-vertex case:

Conjecture

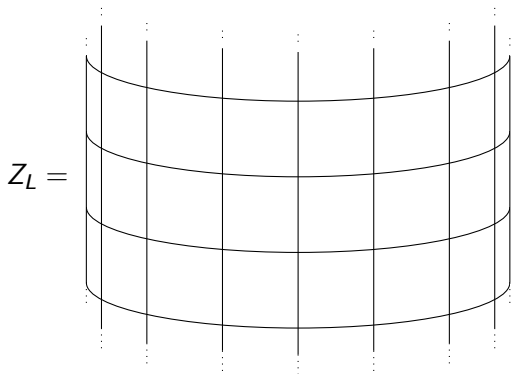
*The normalization of Ψ can be chosen so that its entries are **theta functions** of nome p and of order $L - 1$ in each variable.*

(theta functions = holomorphic, doubly pseudo-periodic functions)

Definition

Define

$$Z_L(x_1, \dots, x_L) = \sum_{\alpha_1, \dots, \alpha_L} \psi_{\alpha_1, \dots, \alpha_L}(x_1, \dots, x_L) \psi_{\alpha_1, \dots, \alpha_L}(-x_1, \dots, -x_L)$$



Characterization

Assuming the conjecture above, we have the

Theorem

Z_L is a symmetric function of its arguments, and a theta function of nome \sqrt{p} and of order $L - 1$ in each.

Furthermore, it is entirely determined by certain recurrence relations (of the same type as Korepin relations for the partition function of the six-vertex model with Domain Wall Boundary Condition).

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Half-specialization

Define

$$Z_L(0, x_1, -x_1, \dots, x_n, -x_n) = \prod_{i=1}^n \vartheta^2(x_i - \eta) \vartheta^2(x_i + \eta) X_n(x_1, \dots, x_n)$$

Then X_n is an even theta function in each of its arguments, of degree $2(2n - 1)$, which is entirely determined by the recurrence relations

- $X_n(\dots, x, x + \eta) = \prod_{i=1}^{n-2} \vartheta(x - \eta - x_i)^4 \vartheta(x - \eta + x_i)^4 \varphi(x) \varphi(x + \eta) X_{n-2}(\dots)$.
- $X_n(\dots, \eta + \alpha) = (\text{stuff}) X_{n-1}(\dots)$ for all three $\alpha \neq 0$ s.t. $2\alpha = 0$.

Theorem

- *The recurrence relations are solved by a product of Pfaffians:*

$$X_n(x_1, \dots, x_{2n}) = A_n(x_1, \dots, x_n) B_n(x_1, \dots, x_n)$$

$$A_n(x_1, \dots, x_n) \propto \text{Pf } f(x_i, x_j), \quad B_n(\dots) = A_{n+1}(\dots, \pi/2 + \eta)$$

where f is a certain skew-symmetric elliptic functions of their arguments. (for A_{2m+1} add extra row/column)

- *A_n and B_n can also be expressed as products of two determinants (similar to elliptic versions of Tsuchiya's determinant as in [Filali, '11]).*
- *In the homogeneous limit, $A_n \rightarrow \zeta^{\lfloor n^2/2 \rfloor} s_n(\zeta^{-2})$ and $B_n \rightarrow (4/3)^n \zeta^{\lfloor (n+1)^2/2 \rfloor} s_{-n-1}(\zeta^{-2})$, and the BM bilinear recurrence relations = Plücker + linear differential relations.*

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Polynomials in the homogeneous limit

If one sends all x_i to zero, then Z_L becomes the product of 4 polynomials in ζ (the constant term corresponding to the six-vertex limit), best expressed separately depending on parity of n :

$$\begin{aligned}
 H_{2m} &= 1, & 1, & 3 + \zeta^2, & 26 + 29\zeta^2 + 8\zeta^4 + \zeta^6 \\
 2^{m-1}H_{2m}(J_2) &= & 1, & 7 + \zeta^2, & 143 + 99\zeta^2 + 13\zeta^4 + \zeta^6 \\
 H_{2m}(J_3) &= & 1, & 2 + \zeta + \zeta^2, & 11 + 12\zeta + 21\zeta^2 + 10\zeta^3 + 7\zeta^4 + \zeta^6 \\
 H_{2m}(J_4) &= & 1, & 2 - \zeta + \zeta^2, & 11 - 12\zeta + 21\zeta^2 - 10\zeta^3 + 7\zeta^4 + \zeta^6 \\
 2^{m-1}H_{2m}(J_2, J_3) &= & 1, & 5 + 2\zeta + \zeta^2, & 66 + 63\zeta + 81\zeta^2 + 30\zeta^3 + 12\zeta^4 + \zeta^6 \\
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 \end{aligned}$$

Polynomials in the homogeneous limit

If one sends all x_i to zero, then Z_L becomes the product of 4 polynomials in ζ (the constant term corresponding to the six-vertex limit), best expressed separately depending on parity of n :

$$H_{2m} = 1, \quad 1, \quad 3 + \zeta^2, \quad 26 + 29\zeta^2 + 8\zeta^4 + \zeta^6$$

1, 1, 3, 26, 646... is the number of VSASMs of odd size.

$$2^{m-1} H_{2m}(J_2) = 1, \quad 7 + \zeta^2, \quad 143 + 99\zeta^2 + 13\zeta^4 + \zeta^6$$

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 \end{aligned}$$

1, 2, 11, 170... is the number of CSTCPPs.

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1, 5, 66, 2431... is one of the factors in the number of UUASMs.

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1, 7, 143, 8398 is ???????

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Vertex-IRF transformation

The equations above do not enable us to compute Ψ because of lack of S^z conservation. \rightarrow use vertex-IRF transformation.

Introduce (dual) **vertex-IRF intertwiners**:

$$t_{\pm}(z)_b^a = \begin{pmatrix} \vartheta_2 \\ \vartheta_3 \end{pmatrix} ((a-b)z + 2a\eta)$$

and for any sequence $(a_i)_{i=0,\dots,L}$ such that $a_{i+1} = a_i \pm 1$,

$$\Phi_{a_0,\dots,a_L} = \sum_{\alpha_1,\dots,\alpha_L=\pm} \Psi_{\alpha_1,\dots,\alpha_L} \prod_{i=1}^L t_{\alpha_i}(z_i)_{a_i}^{a_{i-1}}$$



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Vertex-IRF transformation

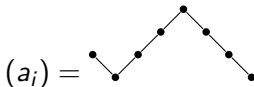
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The exchange relation turns into a simple recurrence formula for Φ .

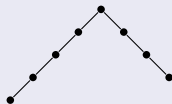
Conjecture (Weston and Z-J, '11)

- $\Phi_{a_0, \dots, a_L} = 0$ if $a_L \neq a_0 \pm 1$.
- The $\Phi_{a, \dots, a+1}$ are determined by the initial condition:

$$\Phi_{a, a+1, \dots, a+n, a+n+1, a+n, \dots, a+1} \propto$$

$$\prod_{1 \leq i < j \leq n+1} \vartheta(z_i - z_j + 2\eta) \prod_{n+2 \leq i < j \leq L} \vartheta(z_i - z_j + 2\eta)$$

$$\prod_{i=n+2}^L \vartheta(z_i) \vartheta_{2,3} \left(\sum_{i=1}^{n+1} z_i - \sum_{i=n+2}^L z_i - 2(a-n)\eta \right)$$



and the recurrence relation: ($\tau_i =$ permutation of z_i and z_{i+1})

$$\Phi_{\dots, b, b-1, b, \dots} = \frac{\vartheta(2\eta) \vartheta(2\eta + z_i - z_{i+1}) \tau_i - \vartheta(2\eta) \vartheta(2b\eta + z_i - z_{i+1})}{\vartheta(2\eta(b+1)) \vartheta(z_{i+1} - z_i)} \Phi_{\dots, b, b+1, b, \dots}$$

Prospects

- Connection to nonsymmetric elliptic Macdonad polynomials?
- Connection/appliation to the work of Rosengren?
- Connection to Painlevé VI?
- Combinatorial interpretation of all the entries of these polynomials?

IRF model:

- Can we “solve” the recurrence relations for the IRF model? (factorized expression?)
- Lashkevich and Pugai introduced elliptic vertex operators and wrote integral expressions for vacuum correlation functions which should be related to Φ_{a_0, \dots, a_L} . Use these integral expressions?
- Can we go back to the “spin” basis and prove other Bazhanov–Mangazeev conjectures?

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