

# Directed polymers and KPZ universality

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(Incomplete) review of directed polymers in i.i.d. random environments, especially KPZ universality in  $1+1$  dimensions

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## 1. General $d + 1$ dimensional model

1.1. Weak and strong disorder.

1.2. Variational formulas, large deviations.

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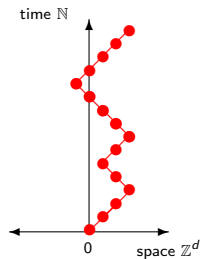
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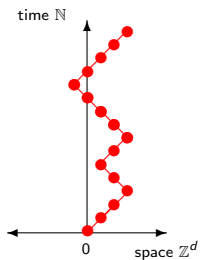
## 2. $1 + 1$ dimensions

- 2.1. KPZ universality
- 2.2. The three exactly solvable models.
- 2.3. Specific results for the **log-gamma polymer**: stationary process, fluctuation exponents, tropical combinatorics.

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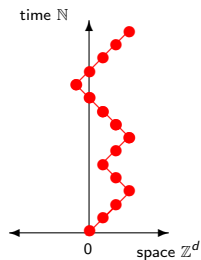


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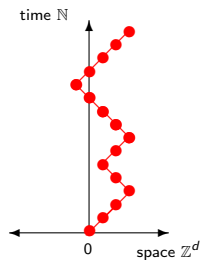
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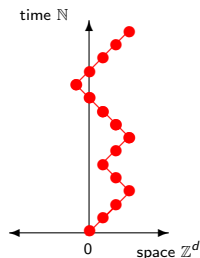
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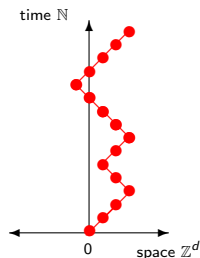
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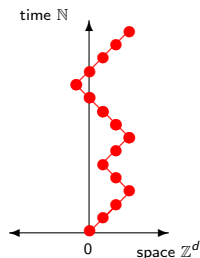
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$\mathbb{P}$  probability distribution on  $\omega$ , often  $\{\omega(k, x)\}$  i.i.d.

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- Dependence on  $\beta$  and  $d$



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**Early rigorous results:** diffusive behavior for  $d \geq 3$  and small  $\beta > 0$ :

1988 Imbrie and Spencer:  $n^{-1}E^Q(|x(n)|^2) \rightarrow c$   $\mathbb{P}$ -a.s.

1989 Bolthausen: quenched CLT for  $n^{-1/2}x(n)$ .

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Kolmogorov's 0-1 law:  $\mathbb{P}(W_\infty > 0) = 0$  or  $1$ .

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**Definition.**  $\left\{ \begin{array}{l} \text{Weak disorder: } W_\infty > 0 \\ \text{Strong disorder: } W_\infty = 0. \end{array} \right.$

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**Theorem.**  $\exists \beta_c \in [0, \infty]$  such that

$\beta \in [0, \beta_c) \implies$  weak disorder

$\beta \in (\beta_c, \infty) \implies$  strong disorder

For  $d \in \{1, 2\}$   $\beta_c = 0$ , while for  $d \geq 3$   $\beta_c \in (0, \infty]$ .

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**Open question:** Are these always the same?

In  $d \in \{1, 2\}$   $\beta_c = \beta'_c = 0$ .

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**Theorem.** Under  $d \geq 3$  and weak disorder,

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**Proof idea.** Construct RWRE  $Q^\omega$  using  $W_\infty > 0$  as a density.

[Comets and Yoshida, 2006]

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## Sufficient conditions for very strong disorder:

- $d = 1$  or  $2$
- $\beta \lambda'(\beta) - \lambda(\beta) > \log(2d)$ . True for some distributions if  $\beta$  large enough.

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$g(\omega, z_{1,\ell})$  is a function on  $\Omega_\ell = \Omega \times \mathcal{R}^\ell$ .

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$$S_z : (\omega, z_{1, \ell}) \mapsto (T_{z_1} \omega, z_{2, \ell} z).$$

Kernel  $p$  on  $\Omega_\ell$ :  $p(\eta, S_z \eta) = |\mathcal{R}|^{-1}$  for  $\eta = (\omega, z_{1, \ell})$ .

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$H_{\mathbb{P}}$  convex but not lower semicontinuous.

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- Analogous result for point-to-point free energy.

Quenched weak LDP (large deviation principle) under  $Q_n$ .

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IID environment, directed walk  $\Rightarrow$  full LDP holds.

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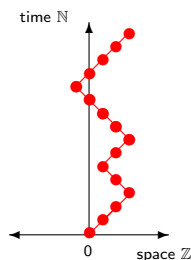
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#### Edwards-Wilkinson (EW)

- time  $\sim n$ , spatial correlations  $\sim n^{1/2}$ , fluctuations  $\sim n^{1/4}$
- Gaussian limits



# KPZ class: 1+1 dim directed polymer



$\{\omega(k, x)\}$  i.i.d. under  $\mathbb{P}$

$$Z_n = E \left[ \exp \left\{ \beta \sum_{k=1}^n \omega(k, X_k) \right\} \right]$$

$$Z_{n,u} = E \left[ \exp \left\{ \beta \sum_{k=1}^n \omega(k, X_k) \right\}, X_n = u \right]$$

$$Q_n(x.) = \frac{1}{Z_n} \exp \left\{ \beta \sum_{k=1}^n \omega(k, x_k) \right\} P(x.)$$

# Expected KPZ behavior

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**Known.**

- Partial results for a handful of exactly solvable models.
- “Weak universality” of Alberts-Khanin-Quastel.



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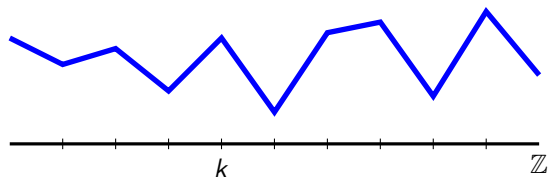
In KPZ class also

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- zero-temperature polymer, or last-passage percolation model
- Other 1+1 dim growth models (PNG, ballistic deposition)
- particle systems with drift and nonlinear flux function (ASEP, ZRP)

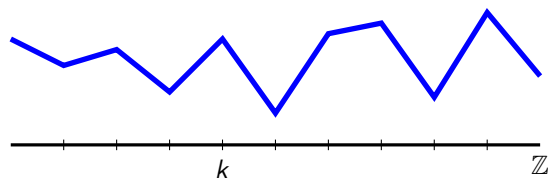
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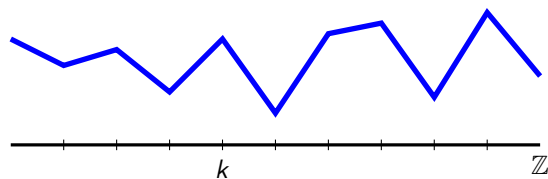


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$\omega_{t,k} = (\omega_{t,k}(j) : |j| \leq R)$  random probability vectors, IID over  $(t, k)$

# RAP scaling limit



$$v = \sum_x x \mathbb{E}\omega(x) \quad \sigma^2 = \sum_x (x - v)^2 \mathbb{E}\omega(x).$$

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## Scaled height process

$$z_n(t, r) = n^{-1/4} \{ \sigma_{\lfloor nt \rfloor}(-\lfloor ntv \rfloor + \lfloor r\sqrt{n} \rfloor) - \mu_0 r \sqrt{n} \}, \quad (t, r) \in \mathbb{R}_+ \times \mathbb{R}.$$

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**Theorem.** [Balázs, Rassoul-Agha, S. 2006]  $z_n(t, r) \Rightarrow Z(t, r)$  where  $Z$  is the Gaussian process

$$Z(t, r) = c_1 \iint_{[0, t] \times \mathbb{R}} \varphi_{\sigma^2(t-s)}(r-x) dW(s, x) + c_2 \int_{\mathbb{R}} \varphi_{\sigma^2 t}(r-x) B(x) dx$$

RAP is an example from the **EW universality class**.

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In this class also

- current of independent random walks (incl. RWRE)
- symmetric simple exclusion process
- Hammersley's serial harness process

## KPZ equation

1986 Kardar, Parisi and Zhang: general model for height function  $h(t, x)$  of a 1+1 dimensional growing interface:

$$h_t = \frac{1}{2} h_{xx} + \frac{1}{2} (h_x)^2 + \dot{W}$$

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**Define**  $h = \log Z$  as the **Hopf-Cole solution** of KPZ.

# KPZ behavior of KPZ equation

- **Balázs, Quastel, and S. (2011):** With initial height function  $h(0, x)$  a two-sided Brownian motion in  $x \in \mathbb{R}$ ,

$$C_1 t^{2/3} \leq \text{Var}(h(t, 0)) \leq C_2 t^{2/3}$$

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**Cross-over distribution** because it has

$$\begin{cases} \text{Tracy-Widom } F_{\text{GUE}} \text{ limit in the scale } t^{1/3} & \text{as } t \nearrow \infty \\ \text{Gaussian limit in the scale } t^{1/4} & \text{as } t \searrow 0. \end{cases}$$

- **Member** of the KPZ universality class because long-term behavior has right exponent and  $F_{\text{GUE}}$  limit.



## Role of KPZ equation

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- Limit of discrete models when asymmetry or noise suitably tuned to zero as the limit is taken.
- First result **Bertini and Giacomin 1997**: height function of **weakly asymmetric** simple exclusion process converges to Hopf-Cole solution of KPZ.

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### 2.2 Exactly solvable directed polymers

Three exactly solvable 1+1 dim models (positive temperature)

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- **Log-gamma polymer** (S 2009).

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Three exactly solvable 1+1 dim models (positive temperature)

- **Continuum directed random polymer**, or Hopf-Cole solution of the KPZ equation, or  $\log Z$  where  $Z$  solves SHE.
- **Semidiscrete polymer**, or cont-time RW paths in Brownian environment (O'Connell-Yor 2001).
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Next brief look at the two discrete models.

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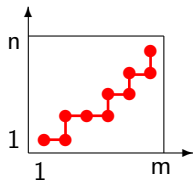
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# Log-gamma polymer

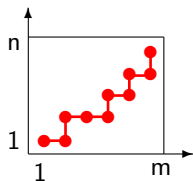


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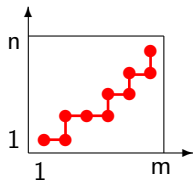
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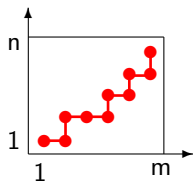
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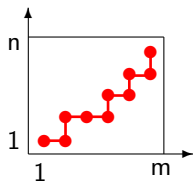
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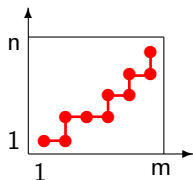
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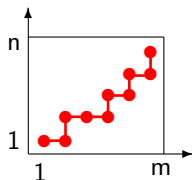
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##### 2. It can be “solved” with ideas from tropical combinatorics

This yields

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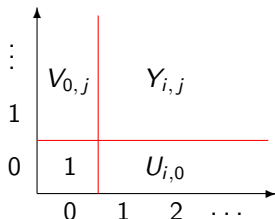
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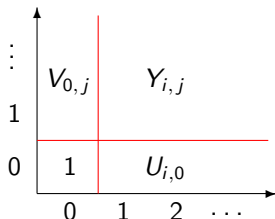
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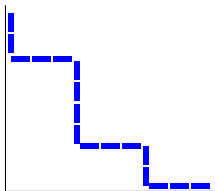


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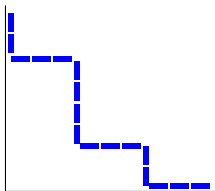


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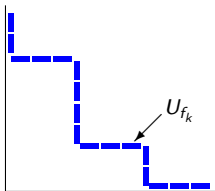
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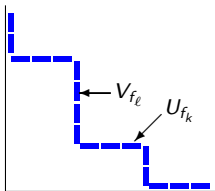
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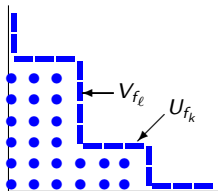
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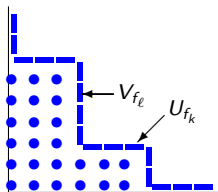


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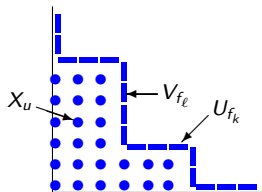
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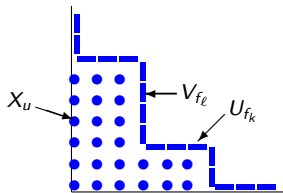
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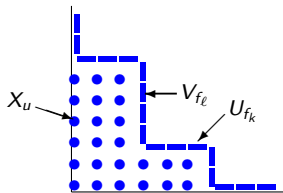
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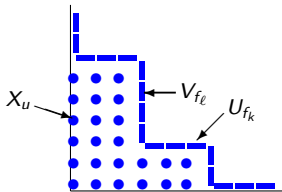
## Theorem.

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## Theorem.

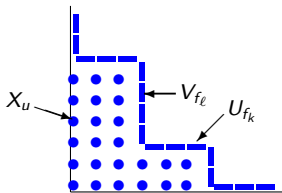
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## Theorem.

For any fixed down-right path,  
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 with marginals

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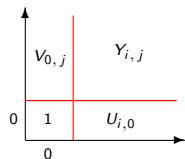
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Hence we could call this the **"Burke property"** of the log-gamma polymer.

# Taking advantage of the stationarity

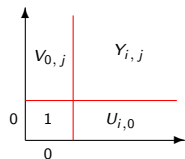


Initial weights ( $i, j \in \mathbb{N}$ ):

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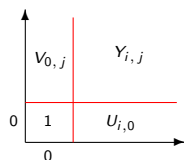
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**Coupling** of two log-gamma models:

- Original one with IID bulk weights, paths  $(1, 1) \rightarrow (m, n)$
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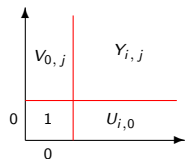
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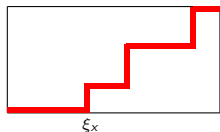
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Let us look at fluctuation exponents for  $\log Z$ .

# Fluctuation exponents: stationary case

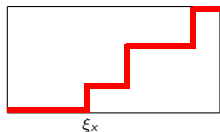
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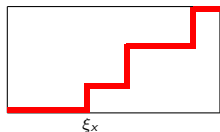
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$$\xi_x = \max\{k \geq 0 : x_i = (i, 0) \text{ for } 0 \leq i \leq k\}$$

For  $\theta, x > 0$  define positive function

$$L(\theta, x) = \int_0^x (\Psi_0(\theta) - \log y) x^{-\theta} y^{\theta-1} e^{x-y} dy$$

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**Theorem.** For the stationary case,

$$\text{Var}[\log Z_{m,n}] = n\Psi_1(\mu - \theta) - m\Psi_1(\theta) + 2 E_{m,n} \left[ \sum_{i=1}^{\xi_x} L(\theta, Y_{i,0}^{-1}) \right]$$

## Remark: polygamma functions

$$\Psi_n(s) = \frac{d^{n+1}}{ds^{n+1}} \log \Gamma(s), \quad n \geq 0$$

These appear naturally because for  $Y^{-1} \sim \text{Gamma}(\mu)$

$$\mathbb{E}(\log Y) = -\Psi_0(\mu) \quad (\text{digamma function})$$

$$\text{Var}(\log Y) = \Psi_1(\mu) \quad (\text{trigamma function})$$

## Fluctuation exponent: stationary case

With  $0 < \theta < \mu$  fixed and  $N \nearrow \infty$  assume

$$|m - N\Psi_1(\mu - \theta)| \leq CN^{2/3} \quad \text{and} \quad |n - N\Psi_1(\theta)| \leq CN^{2/3} \quad (1)$$



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Theorem: Off-characteristic CLT

Suppose  $n = \Psi_1(\theta)N$  and  $m = \Psi_1(\mu - \theta)N + \gamma N^\alpha$  with  $\gamma > 0$ ,  $\alpha > 2/3$ .  
Then

$$N^{-\alpha/2} \left\{ \log Z_{m,n} - \mathbb{E}(\log Z_{m,n}) \right\} \Rightarrow \mathcal{N}(0, \gamma \Psi_1(\theta))$$

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$$p_{s,t}(\mu) \equiv \lim_{N \rightarrow \infty} \frac{\log Z_{Ns, Nt}}{N} = \inf_{\theta \in (0, \mu)} \{-s\Psi_0(\theta) - t\Psi_0(\mu - \theta)\}$$

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**Theorem.** Upper bound for fluctuation exponent:

$$\mathbb{P}\left\{ \left| \log Z_{Ns, Nt} - N\rho_{s,t}(\mu) \right| \geq bN^{1/3} \right\} \leq Cb^{-3/2}$$

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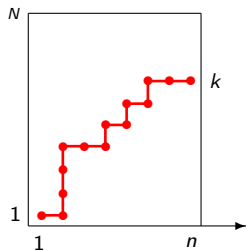
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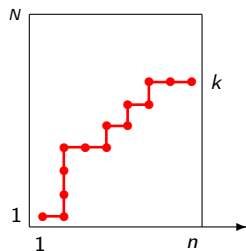
**Remark.** Corresponding bounds exist for path with KPZ exponent  $2/3$ .

# Combinatorial approach to log-gamma polymer





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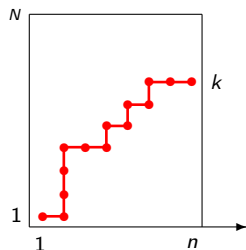


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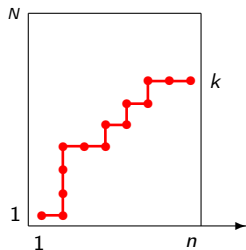
$$\Pi_{n,k}^1 = \{ \text{admissible paths } (1, 1) \rightarrow (n, k) \}$$

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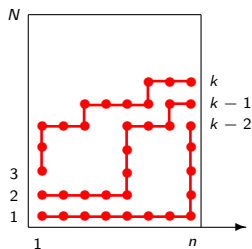


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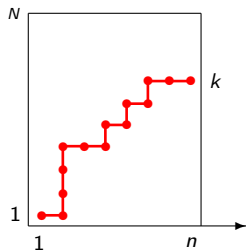
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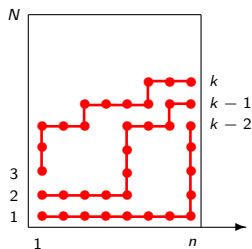


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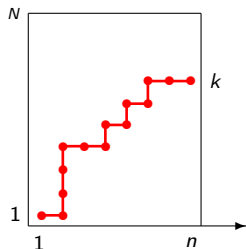
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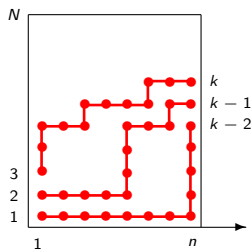


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- Details not illuminating.



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Theory of **Markov functions** is useful here.

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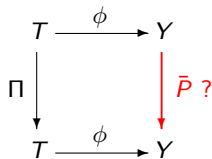
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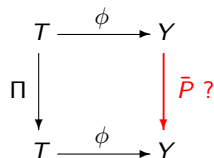


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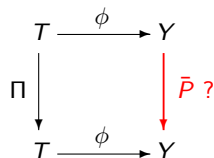
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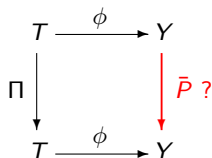
Set  $w(y) = K(y, T)$ .

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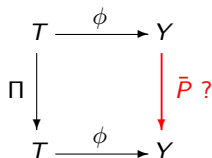
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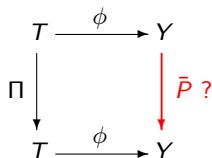
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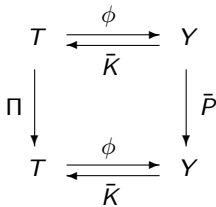
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Then  $\bar{K} \circ \Pi = \bar{P} \circ \bar{K}$

$$\begin{array}{ccc} T & \begin{array}{c} \xrightarrow{\phi} \\ \xleftarrow{\bar{K}} \end{array} & Y \\ \Pi \downarrow & & \downarrow \bar{P} \\ T & \begin{array}{c} \xrightarrow{\phi} \\ \xleftarrow{\bar{K}} \end{array} & Y \end{array}$$



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**Theorem.** [Rogers and Pitman, 1981]

If  $z(n)$  starts with distribution  $\bar{K}(y, dz)$ , then  $y(n)$  is Markov in its own filtration with transition  $\bar{P}$  and initial state  $y(0) = y$ .

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and intertwining kernel  $K : \mathbb{Y}_N \rightarrow \mathbb{T}_N$  by

$$K(y, dz) = \prod_{1 \leq l \leq k < N} \left(\frac{z_{k,l}}{z_{k+1,l}}\right)^{\theta_{k+1} - \theta_l} \\ \times \exp\left(-\frac{z_{k,l}}{z_{k+1,l}} - \frac{z_{k+1,l+1}}{z_{k,l}}\right) \frac{dz_{k,l}}{z_{k,l}} \prod_{l=1}^N \delta_{y_l}(dz_{N,l})$$

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Beneficial because known special functions diagonalize the transition kernel.



# Whittaker functions

$GL(N, \mathbb{R})$ -**Whittaker function** is given for  $y \in \mathbb{Y}_N$ , with  $\lambda \in \mathbb{C}^N$ , by

$$\Psi_\lambda(y) = \prod_{i=1}^N y_i^{-\lambda_i} \int_{\mathbb{T}_N} K_\lambda(y, dz)$$

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Intertwining develops into

$$\int_{(0, \infty)^N} \frac{\Psi_{\theta+\lambda}(\tilde{y})}{\Psi_\theta(\tilde{y})} \bar{P}_n(y, d\tilde{y}) = \left( \prod_{i=1}^N \frac{\Gamma(\gamma_{n,i} + \lambda_i)}{\Gamma(\gamma_{n,i})} \right) \frac{\Psi_{\theta+\lambda}(y)}{\Psi_\theta(y)}$$

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**Future goal:** asymptotics for distribution of  $\log z_{N,1}(n)$ ?

# Work in progress: intermediate disorder exponents

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**Theorem.** These exponents valid for stationary semidiscrete polymer.  
Upper bounds valid for model without boundaries. [Moreno, S, Valkó]

# Explicit large deviations for $\log Z$

L.m.g.f. of  $\log Y$ ,  $Y \sim \Gamma^{-1}(\mu)$ :

$$M_{\mu}(\xi) = \log \mathbb{E}(e^{\xi \log Y}) = \begin{cases} \log \Gamma(\mu - \xi) - \log \Gamma(\mu) & \xi \in (-\infty, \mu) \\ \infty & \xi \in [\mu, \infty). \end{cases}$$

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**Theorem.** [Georgiou, S 2011]

$$\Lambda_{s,t}(\xi) = \begin{cases} p(s,t)\xi & \xi < 0 \\ \inf_{\theta \in (\xi, \mu)} \{tM_{\theta}(\xi) - sM_{\mu-\theta}(-\xi)\} & 0 \leq \xi < \mu \\ \infty & \xi \geq \mu. \end{cases}$$

- $\Lambda_{s,t}$  linear on  $\mathbb{R}_-$  because for  $r < p(s, t)$

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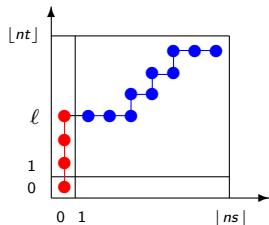
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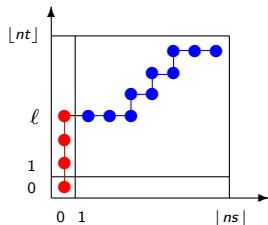
- Proof of formula for  $\Lambda_{s,t}$  goes by first finding  $J_{s,t}$  and then convex conjugation.

# Starting point for proof of large deviations



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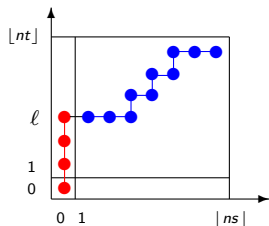


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Divide by  $\prod_{j=1}^{nt} V_{0, j}$  :

$$\prod_{i=1}^{ns} U_{i, nt} = \sum_{\ell=1}^{nt} \left( \prod_{j=\ell+1}^{nt} V_{0, j}^{-1} \right) Z_{(1, \ell), (ns, nt)} \\ + \sum_{k=1}^{ns} \left( \prod_{j=1}^{nt} V_{0, j}^{-1} \right) \left( \prod_{i=1}^k U_{i, 0} \right) Z_{(k, 1), (ns, nt)}$$

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Now we know LDP for  $\log(\text{l.h.s})$  and can extract  $\log Z$  from the r.h.s.