# 2013 Summer Graduate Workshop, Cortona, Italy: <br> Mathematical General Relativity <br> Some Scalar Curvature Basics 

Recall that for better or worse, we defined the curvature tensor as

$$
R(X, Y, Z)=\nabla_{X} \nabla_{Y} Z-\nabla_{Y} \nabla_{X} Z-\nabla_{[X, Y]} Z,
$$

with $R\left(\frac{\partial}{\partial x^{i}}, \frac{\partial}{\partial x^{j}}, \frac{\partial}{\partial x^{k}}\right)=R_{i j k}^{\ell} \frac{\partial}{\partial x^{\ell}}$.
Also, recall that a comma in a subscript denotes partial differentiation (with respect to some coordinate), whereas a semicolon in a subscript denotes covariant differentiation. For example, if $h$ is a (1,2)-tensor with components in a coordinate chart $h_{j k}^{i}$, then the covariant derivative $\nabla h$ is a ( 1,3 )-tensor with components
$h_{j k ; \ell}^{i}:=\nabla h\left(d x^{i}, \frac{\partial}{\partial x^{j}}, \frac{\partial}{\partial x^{k}}, \frac{\partial}{\partial x^{\ell}}\right)=\left(\nabla_{\frac{\partial}{\partial x^{\ell}}} h\right)\left(d x^{i}, \frac{\partial}{\partial x^{j}}, \frac{\partial}{\partial x^{k}}\right)=h_{j k, \ell}^{i}+\Gamma_{m \ell}^{i} h_{j k}^{m}-\Gamma_{j \ell}^{m} h_{m k}^{i}-\Gamma_{k \ell}^{m} h_{j m}^{i}$.
Problem 1. Linearization of Scalar Curvature. Let $R(g)=g^{i j} R_{i j}$ be the scalar curvature of a metric (not necessarily Riemannian). Consider a variation $g(t)=g+t h$ of $g$ in the direction of a symmetric ( 0,2 )-tensor $h$ (more generally, note that all you will use is that $g(t)$ is smooth in $t$, with $g(0)=g$ and $g^{\prime}(0)=h$. For small $t, g(t)$ is a metric. Define $L_{g}(h):=D R_{g}(h)=\left.\frac{d}{d t}\right|_{t=0} R(g(t))$.

Derive the scalar curvature formula

$$
R(g)=g^{i j} R_{i j}=g^{i j}\left(\Gamma_{i j, k}^{k}-\Gamma_{i k, j}^{k}+\Gamma_{k \ell}^{k} \Gamma_{i j}^{\ell}-\Gamma_{j \ell}^{k} \Gamma_{i k}^{\ell}\right)
$$

and use it to verify the identity

$$
L_{g}(h)=-\Delta_{g}\left(\operatorname{tr}_{g}(h)\right)+\operatorname{div}_{g}\left(\operatorname{div}_{g}(h)\right)-\langle h, \operatorname{Ric}(g)\rangle_{g} .
$$

Hints: One approach is to compute in normal coordinates at a point $p$, so that $g_{i j}(p)=\delta_{i j}$, as well as $\Gamma_{i j}^{k}(p)=0$, equivalently, $g_{i j, k}(p)=0$. Indicated below are some formulas to guide you, and which you should verify as you compute. We emphasize here and below that all quantities are evaluated at $p$. In such normal coordinates, we have at $p: h_{i j ; k \ell}=h_{i j, k \ell}-h_{m j} \Gamma_{i k, \ell}^{m}-h_{i m} \Gamma_{j k, \ell}^{m}$. From this, one shows that at $p$,

$$
\begin{aligned}
&-\Delta_{g}\left(\operatorname{tr}_{g}(h)\right)=-g^{k \ell} g^{i j} h_{i j ; k \ell}=-\sum_{i, k}\left(h_{i i, k k}-2 h_{k m} \Gamma_{i k, i}^{m}\right) \\
& \operatorname{div}_{g}\left(\operatorname{div}_{g}(h)\right)=g^{j \ell} g^{i k} h_{i j ; k \ell}=\sum_{i, k}\left(h_{i k, i k}-h_{k m} \Gamma_{i i, k}^{m}-h_{k m} \Gamma_{i k, i}^{m}\right) .
\end{aligned}
$$

To find the variation of the scalar curvature, express the scalar curvature in terms of Christoffel symbols, and take a time derivative, expand and turn the crank. To find the derivative of the inverse metric, note that if $A(t)$ is a smooth curve in $G L(n)$, then $\frac{d}{d t} A^{-1}(t)$ is easily computed from $A(t) A^{-1}(t)=I$ using the product rule.

Another approach is to note that $\left.\frac{d}{d t}\right|_{t=0} \Gamma_{i j}^{k}$ are the components $(\delta \Gamma)_{i j}^{k}$ of a tensor $\delta \Gamma$ (this follows from \#4 a on the earlier set on basic geometry problems), so that clearly the variation of the Ricci tensor is given by $\left.\frac{d}{d t}\right|_{t=0} R_{i j}=(\delta \Gamma)_{i j ; k}^{k}-(\delta \Gamma)_{i k ; j}^{k}$. One should now express $\delta \Gamma$ in terms of the covariant derivative of $h$.

## Problem 2. Conformal deformation of scalar curvature.

a. Suppose $\left(M^{n}, g\right)$ is a Riemannian metric and $\tilde{g}=e^{\varphi} g$. Show that

$$
R(\tilde{g})=e^{-\varphi}\left(R(g)-(n-1) \Delta_{g} \varphi-\frac{1}{4}(n-1)(n-2)|\nabla \varphi|_{g}^{2}\right) .
$$

b. In case $n \geq 3$, if we write $e^{\varphi}=u^{\frac{4}{n-2}}$ for $u>0$, then

$$
R(\tilde{g})=u^{-\frac{n+2}{n-2}}\left(R(g) u-\frac{4(n-1)}{(n-2)} \Delta_{g} u\right)
$$

c. Let $c(n)=\frac{n-2}{4(n-1)}$ and $L_{g} u=\Delta_{g} u-c(n) R(g) u$ is the conformal Laplacian, show that the total scalar curvature of $\tilde{g}=u^{\frac{4}{n-2}} g$ is given by

$$
\int_{M} R(\tilde{g}) d v_{\tilde{g}}=c(n)^{-1} \int_{M}\left(|\nabla u|_{g}^{2}+c(n) R(g) u^{2}\right) d v_{g}
$$

Hint: Show that $d v_{\tilde{g}}=u^{\frac{2 n}{n-2}} d v_{g}$.

## Problem 3. Volume expansion of geodesic balls.

a. Suppose $(V,\langle\rangle$,$) is an n$-dimensional real inner product space. Suppose that $T: V \rightarrow V$ is a self-adjoint linear operator. If $d \sigma$ is the Euclidean area measure, $\mathbb{B}^{n}$ is the unit ball, and $\mathbb{S}^{n-1} \subset V$ is the unit sphere in $V$, then if vol is the Euclidean volume,

$$
\int_{x \in \mathbb{S}^{n-1}}\langle T(x), x\rangle d \sigma=\operatorname{trace}(T) \operatorname{vol}\left(\mathbb{B}^{n}\right) .
$$

b. If $(M, g)$ is Riemannian and $p \in M$, let $B_{r}(p) \subset M$ be the geodesic ball of radius $r>0$ (for sufficiently small $r$ ). Then

$$
\operatorname{vol}_{g}\left(B_{r}(p)\right)=\operatorname{vol}\left(\mathbb{B}^{n}\right) r^{n}\left[1-\frac{\left.R(g)\right|_{p}}{6(n+2)} r^{2}+O\left(r^{3}\right)\right] .
$$

Hint: You may wish to observe and use $\operatorname{det}(I+t A)=1+t \operatorname{trace}(A)+O\left(t^{2}\right)$, along with Problem \#4c below.

Problem 4. Metric expansion in normal coordinates. Suppose that $\gamma(t)$ is a unit-speed geodesic, and that $J(t)$ is a Jacobi field along $\gamma: J^{\prime \prime}(t)=R\left(\gamma^{\prime}(t), J(t), \gamma^{\prime}(t)\right)$.
a. If $R$ is the Riemann curvature tensor, show that

$$
J^{\prime \prime \prime}(t)=\left(\nabla_{\gamma^{\prime}(t)} R\right)\left(\gamma^{\prime}(t), J(t), \gamma^{\prime}(t)\right)+R\left(\gamma^{\prime}(t), J^{\prime}(t), \gamma^{\prime}(t)\right) .
$$

b. Suppose $J(0)=0$. Let $\chi(t)=\langle J(t), J(t)\rangle$. Derive the fourth-order Taylor expansion

$$
\chi(t)=\sum_{k=0}^{4} \frac{\chi^{(k)}(0)}{k!} t^{k}+\mathcal{E}(t)=\left|J^{\prime}(0)\right|^{2} t^{2}-\frac{1}{3}\left\langle R\left(\gamma^{\prime}(0), J^{\prime}(0), J^{\prime}(0), \gamma^{\prime}(0)\right)\right\rangle t^{4}+O\left(t^{5}\right) .
$$

c. Suppose $(M, g)$ is a Riemannian manifold and $p \in M$. Show that in normal coordinates centered at $p$ (so $\left.x^{i}(p)=0\right)$

$$
g_{i j}(x)=\delta_{i j}-\frac{1}{3} R_{k i j \ell} x^{k} x^{\ell}+O\left(|x|^{3}\right) .
$$

Hint: In normal coordinates $\left(x^{i}\right)$ centered at $p$, consider a unit speed radial geodesic $\gamma(t)$ and the vector field $W(t)=t W^{i} \frac{\partial}{\partial x^{i}}$ along $\gamma$, where $W^{i}$ are constants. Show that $W(t)$ is a Jacobi field along $\gamma$ with $W^{\prime}(0)=\left.W^{i} \frac{\partial}{\partial x^{i}}\right|_{p}$. One way to do this is to build a variation $\Gamma(s, t)$ of $\gamma=\Gamma(0, \cdot)$ through geodesics. In any case, you might first observe that in normal coordinates, the curve $\beta$ with components $\beta^{i}(t)=t V^{i}$ is a geodesic with $\beta(0)=p$ and $\beta^{\prime}(0)=\left.V^{i} \frac{\partial}{\partial x^{i}}\right|_{p}$.
Problem 5. Geometric formula for Gaussian curvature. Let $(M, g)$ be a surface with a Riemannian metric $g$. Consider an orthonormal basis $\mathbf{e}_{1}, \mathbf{e}_{2}$ of $T_{p} M$. Note that the Gauss curvature at $p$ is just $K(p)=R\left(\mathbf{e}_{1}, \mathbf{e}_{2}, \mathbf{e}_{2}, \mathbf{e}_{1}\right)$, where $R$ is the Riemann tensor. Consider a normal neighborhood of radius $a>0$ about $p$, with normal coordinates $(x, y)$ built off of the orthonormal basis $\mathbf{e}_{1}, \mathbf{e}_{2}$ of $T_{p} M: \quad(x, y) \mapsto \exp _{p}\left(x \mathbf{e}_{1}+y \mathbf{e}_{2}\right)$. Define geodesic polar coordinates by $(r, \theta) \mapsto f(r, \theta)=\exp _{p}\left(r \cos \theta \mathbf{e}_{1}+r \sin \theta \mathbf{e}_{2}\right)$. Note that the change of coordinates map is just $(r, \theta) \mapsto(x, y)=(r \cos \theta, r \sin \theta)$, which shows that the map $f$, which is clearly smooth for $r<a$ and all $\theta$, is a diffeomorphism for $0<r<a$ and $\theta \in I$, where $I$ in any open interval of length at most $2 \pi$. Note that by the Gauss' Lemma, the metric components in geodesic polar coordinates are $g_{r r}=1, g_{r \theta}=0$ and $g_{\theta \theta}=\left|\frac{\partial f}{\partial \theta}\right|^{2}$. Since the radial curves of constant $\theta$ on $M$ are geodesics, for any $\theta_{0}, J(r)=\frac{\partial f}{\partial \theta}\left(r, \theta_{0}\right)$ is a Jacobi field along the radial geodesic $r \mapsto \gamma(r)=f\left(r, \theta_{0}\right)$.
a. For any $\theta_{0}$, show that $J^{\prime}(0) \perp \gamma^{\prime}(0)$.
b. Use Problem 4 to show that $g_{\theta \theta}(r, \theta)=r^{2}-\frac{K(p)}{3} r^{4}+\mathcal{E}(r, \theta)$, where $\mathcal{E}(r, \theta)=O\left(r^{5}\right)$ uniformly in $\theta$, i.e. $R(r, \theta) \leq C r^{5}, C$ can be chosen independent of $r$ and $\theta$. Use the Taylor expansion $(1+x)^{\alpha}=1+\alpha x+O\left(x^{2}\right)$ to derive

$$
\sqrt{g_{\theta \theta}}(r, \theta)=r-\frac{K(p)}{3!} r^{3}+O\left(r^{4}\right) .
$$

c. Let $L(r)$ be the length of a geodesic circle of radius $r$ about $p$, and let $A(r)$ be the area enclosed by this circle, both computed using the metric $g$. Show that

$$
\lim _{r \rightarrow 0^{+}} \frac{3}{\pi} \frac{2 \pi r-L(r)}{r^{3}}=K(p)=\lim _{r \rightarrow 0^{+}} \frac{12}{\pi} \frac{\pi r^{2}-A(r)}{r^{4}} .
$$

d. Let $D$ be the unit disk in the plane: $D=\left\{(x, y): x^{2}+y^{2}<1\right\}$. Consider the hyperbolic metric

$$
g_{H}=\frac{4}{\left(1-\left(x^{2}+y^{2}\right)\right)^{2}}\left(d x^{2}+d y^{2}\right)=\frac{4}{\left(1-\rho^{2}\right)^{2}}\left(d \rho^{2}+\rho^{2} d \theta^{2}\right)
$$

where $(\rho, \theta)$ are standard polar coordinates on $D$. By solving the differential equation $\frac{2 d \rho}{1-\rho^{2}}=d r$, show how to re-write the hyperbolic metric as $g_{H}=d r^{2}+\sinh ^{2} r d \theta^{2}$. Use this along with the formulas above to show $K=-1$ at the origin of coordinates (of course, $K=-1$ everywhere). (Recall that $d r^{2}=d r \otimes d r$.)

