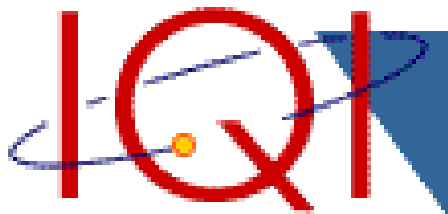


Entanglement in quantum critical phenomena

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Entanglement in quantum critical phenomena

Joint work with

Jose Ignacio Latorre,

Enrique Rico

University of Barcelona

and

Alexei Kitaev

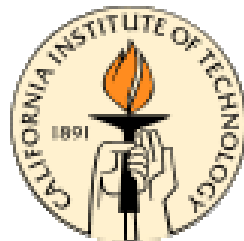
IQI, Caltech



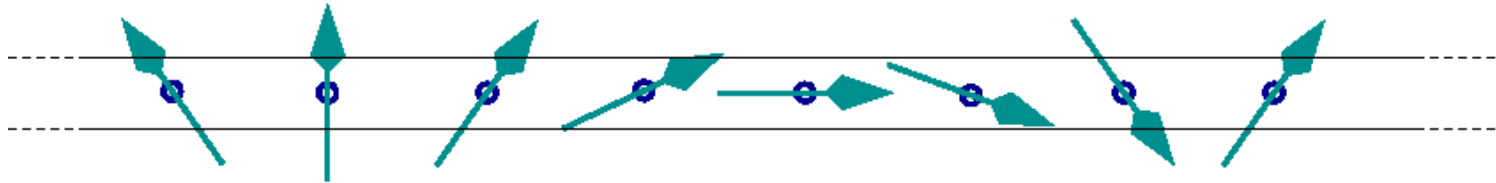
Overview

Vidal, Latorre, Rico, Kitaev,
quant-ph/0211074

1. 1D Spin models and quantum phase transitions
2. Entanglement in spin chains
 - *Definitions*
 - *Computation*
3. Critical and non-critical entanglement
 - Breakdown of DMRG techniques
4. Connection with conformal field theory
 - *Entanglement in 2D and 3D spin models*
 - Monotonicity of entanglement along RG flow



1D spin models

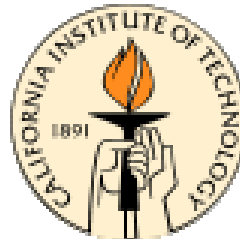


- ***XY model*** with magnetic field
[including ***XX model*** and ***Ising model***]

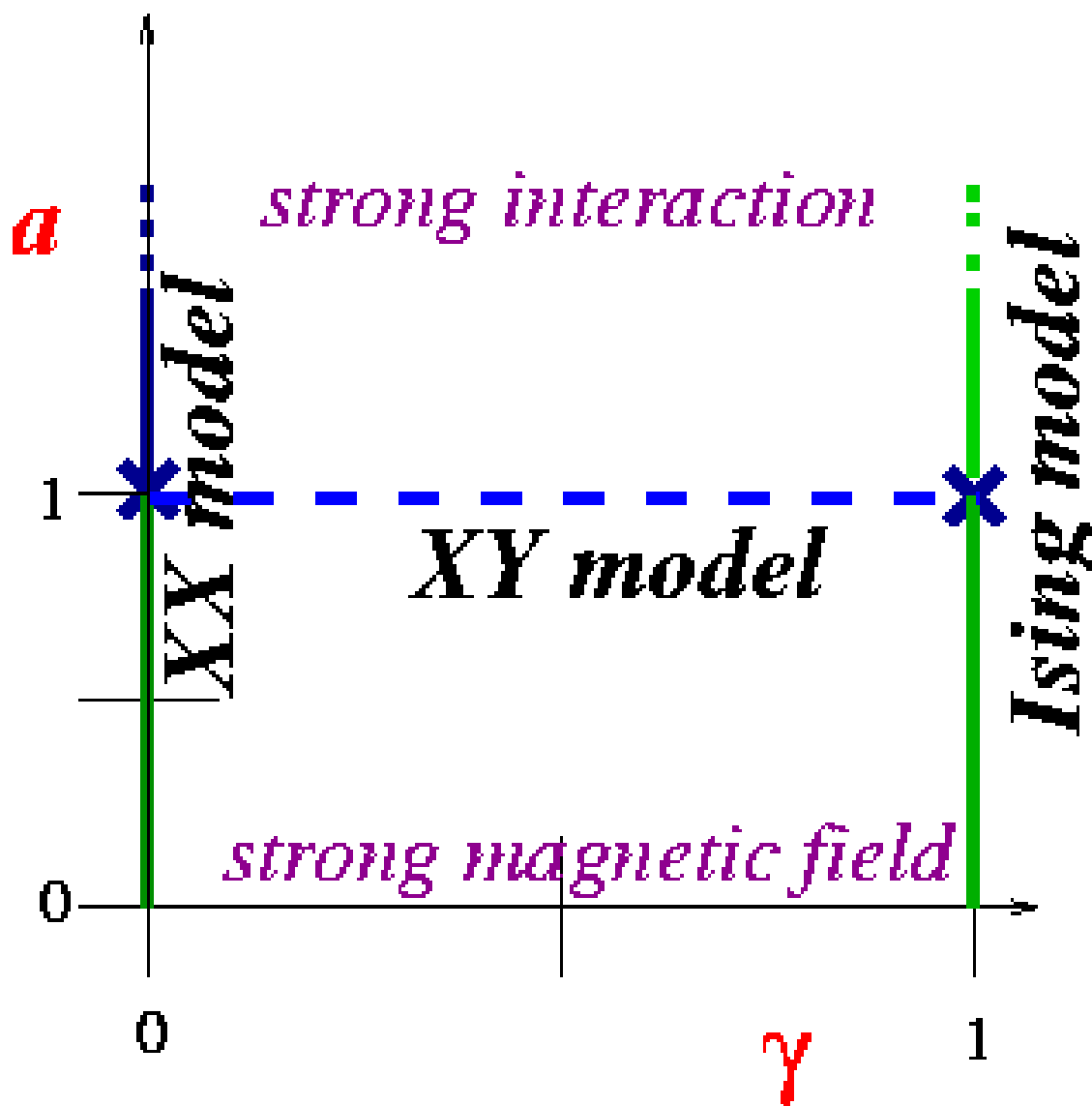
$$H_{XY} = -\sum_{l=0}^{N-1} \left(\frac{a}{2} \left[(1+\gamma)\sigma_l^x \sigma_{l+1}^x + (1-\gamma)\sigma_l^y \sigma_{l+1}^y \right] + \sigma_l^z \right)$$

- ***XXZ model*** with magnetic field

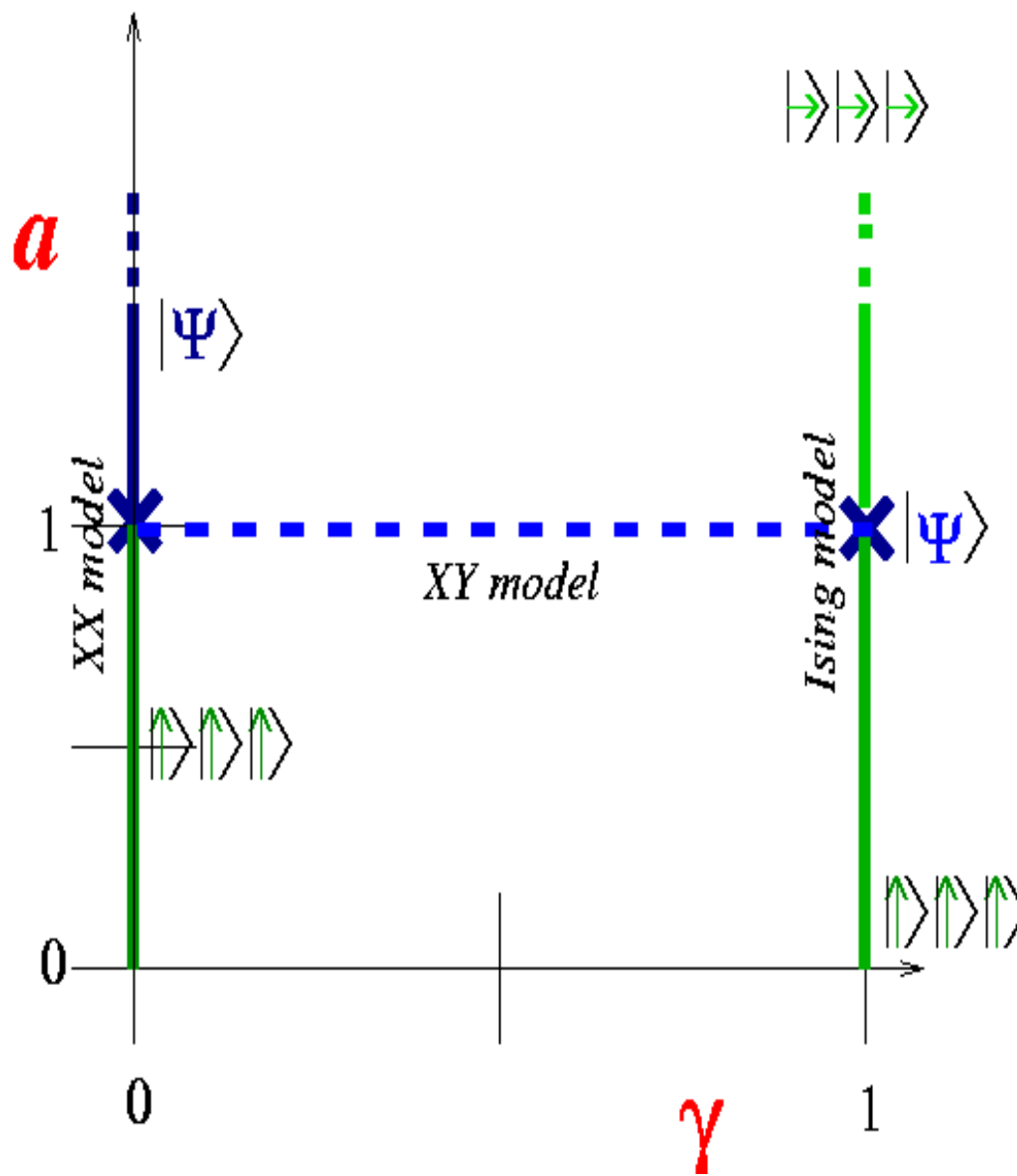
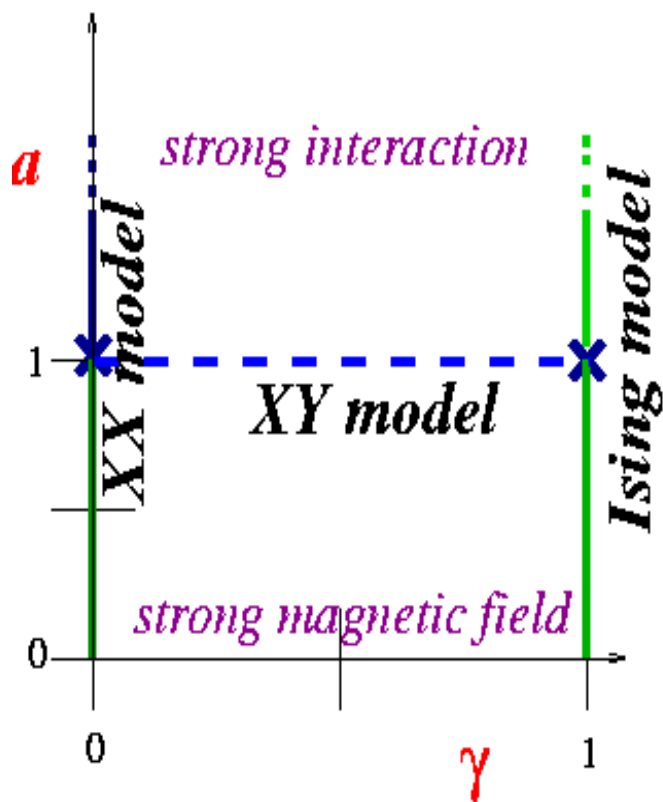
$$H_{XY} = -\sum_{l=0}^{N-1} \left(\sigma_l^x \sigma_{l+1}^x + \sigma_l^y \sigma_{l+1}^y + \Delta \sigma_l^z \sigma_{l+1}^z + \lambda \sigma_l^z \right)$$



XY model



Entanglement and phases



Quantum phase transition

- $T=0$
- $H = H_0 + g H_1$
non-analyticity of ground-state energy
 - Level crossing, $[H_0, H_1] \neq 0$ (finite chain)
 - Thermodynamic limit (infinite chain)
- Long-range correlations

$$\langle \sigma_l \sigma_{l+d} \rangle \propto \frac{1}{d^p}$$



Previous work

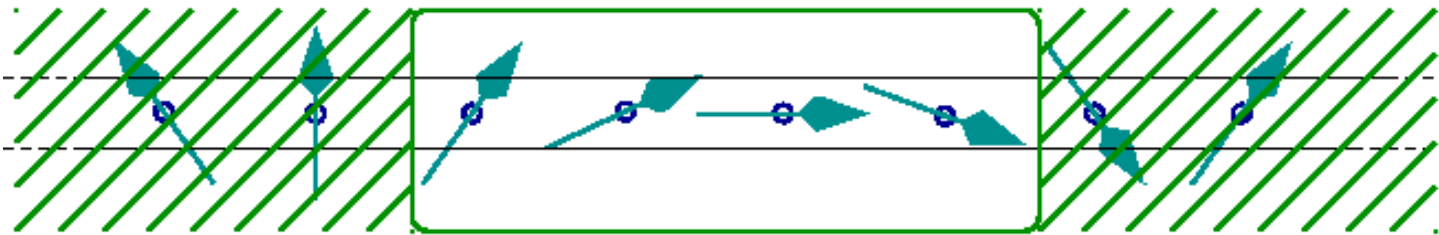
- Dorit Aharonov, Phys. Rev. A 62 (2000);
Transition from quantum to classical in noisy quantum computers.
- Osborne and Nielsen [quant-ph/0109024; quant-ph/0202162]
- Osterloh, Amico, Falci and Falzio, Nature 416 (2002)

Single spin and two-spin entanglement measures have a peak at or close to a phase transition



Entanglement in a spin chain

- We measure the entanglement between a block **B** of spins and the rest of the chain



$$\rho_B = \text{tr}_{\text{chain}-B} \left| \Psi_g \right\rangle \left\langle \Psi_g \right|$$

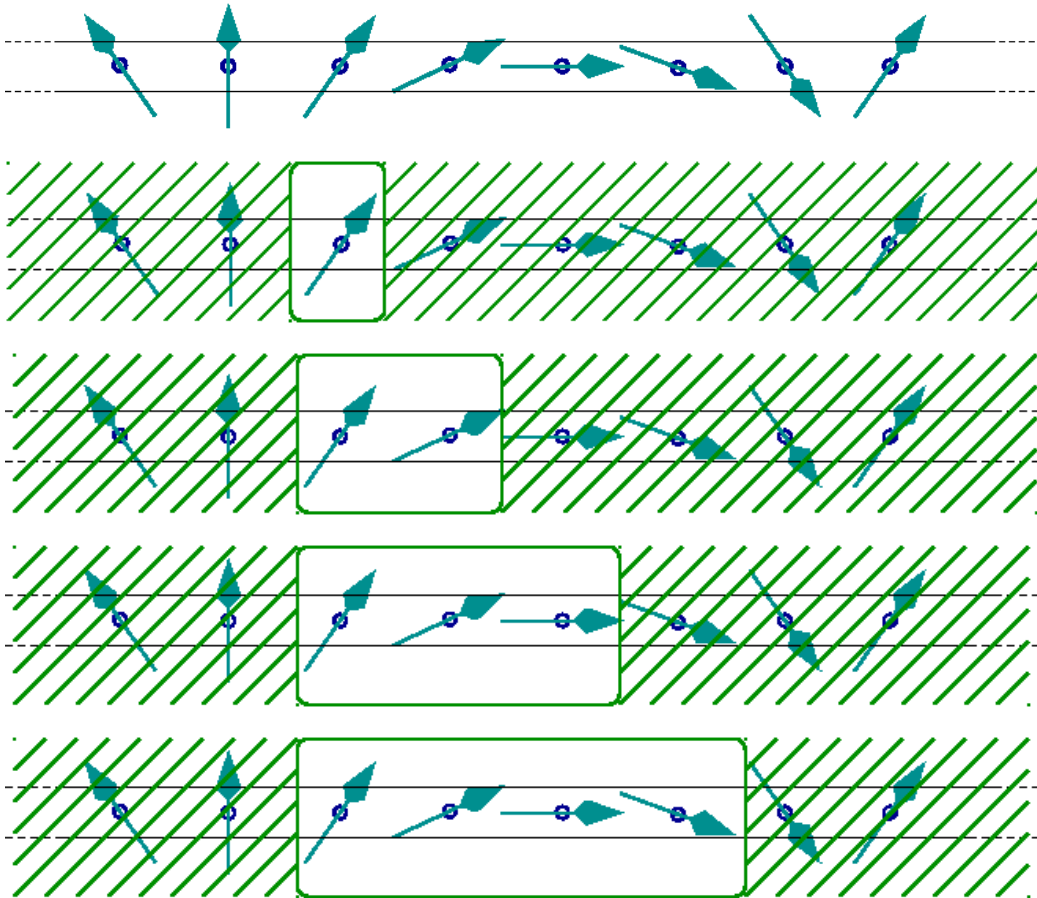
Entropy of entanglement

$$S_B = -\text{tr}(\rho_B \log \rho_B)$$

Bennett, Bernstein, Popescu and
Schumacher, PRA 53 (1996)



Entropy of entanglement



Osborne and Nielsen,
quant-ph/0202162

$$\rho_1 \rightarrow S_1 = S(\rho_1)$$

$$\rho_2 \rightarrow S_2 = S(\rho_2)$$

$$\rho_3 \rightarrow S_3 = S(\rho_3)$$

$$\rho_4 \rightarrow S_4 = S(\rho_4)$$

.....

Computation of entanglement: infinite XY spin chain

- Change of variables

$$c_{2l} = \left(\prod_{m=0}^{l-1} \sigma_m^z \right) \sigma_l^x \quad c_{2l+1} = \left(\prod_{m=0}^{l-1} \sigma_m^z \right) \sigma_l^y$$

Majorana operators

$$c_m^+ = c_m$$

$$\{c_m, c_n\} = 2\delta_{mn}$$

$$H_{XY}(\{\sigma_l^\alpha\}) \rightarrow H_{XY}(\{c_l\})$$

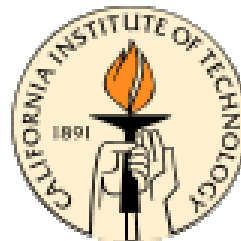


- The ground state $|\Psi_g\rangle$ is **gaussian**
- Compute **correlation matrix**

$$\langle c_m c_n \rangle = \delta_{mn} + iB_{mn} \quad m, n = 0, 1, \dots, 2N-1$$

$$B = \begin{bmatrix} \Pi_0 & \Pi_1 & \cdots & \Pi_{N-1} \\ \Pi_{-1} & \Pi_0 & & \vdots \\ \vdots & & \ddots & \vdots \\ \Pi_{1-N} & \cdots & \cdots & \Pi_0 \end{bmatrix} \quad \Pi_0 = \begin{bmatrix} 0 & g_l \\ -g_{-l} & 0 \end{bmatrix}$$

$$g_l = \frac{1}{2\pi} \int_0^{2\pi} d\phi \exp(-il\phi) \frac{a \cos \phi - 1 - ia\gamma \sin \phi}{|a \cos \phi - 1 - ia\gamma \sin \phi|}$$



- The state ρ_L is also **gaussian**
- Compute **correlation matrix**

$$\langle c_m c_n \rangle = \delta_{mn} + i(B_L)_{mn} \quad m, n = 0, 1, \dots, 2L-1$$

$$B = \begin{bmatrix} \Pi_0 & \Pi_1 & \cdots & \cdots & \Pi_{N-1} \\ \Pi_{-1} & \Pi_0 & & & \vdots \\ \vdots & & \ddots & & \vdots \\ \vdots & & & \ddots & \vdots \\ \Pi_{1-N} & \cdots & \cdots & \cdots & \Pi_0 \end{bmatrix} \rightarrow B_L = \begin{bmatrix} \Pi_0 & \Pi_1 & \cdots & \Pi_{L-1} \\ \Pi_{-1} & \Pi_0 & & \vdots \\ \vdots & & \ddots & \vdots \\ \Pi_{1-L} & \cdots & \cdots & \Pi_0 \end{bmatrix}$$



- Diagonalize *correlation matrix*

$$B_L' = VB_LV^+ = \begin{bmatrix} v_1\Pi & & & \\ & v_2\Pi & & \\ & & \ddots & \\ & & & v_L\Pi \end{bmatrix} \quad \Pi = \begin{bmatrix} 0 & 1 \\ -1 & 0 \end{bmatrix}$$

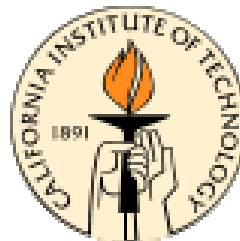
New majorana operators

$$\{c_m\} \xrightarrow{V} \{d_m\}$$

$$d_m^+ = d_m$$

$$\{d_m, d_n\} = 2\delta_{mn}$$

$$\langle d_m d_n \rangle = \delta_{mn} + i(B_L')_{mn}$$



- Introduce *fermionic operators*

$$a_l = \frac{1}{2}(d_{2l} + id_{2l+1})$$

$$\{a_m, a_n\} = \{a_m^+, a_n^+\} = 0$$

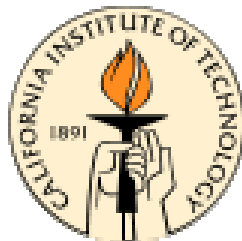
$$\{a_m^+, a_n\} = \delta_{mn}$$

Their *correlation matrix* is $\langle b_m b_n \rangle = 0$

$$\langle b_m^+ b_n \rangle = \delta_{mn} \frac{1 + \nu_m}{2}$$

Therefore the L fermionic modes are *uncorrelated*

$$\rho_L = \sigma_1 \otimes \sigma_2 \otimes \dots \otimes \sigma_L$$



- Therefore the entropy of mode m is

$$S(\sigma_m) = H_2\left(\frac{1+v_m}{2}\right)$$

- And the entropy of the L spins reads

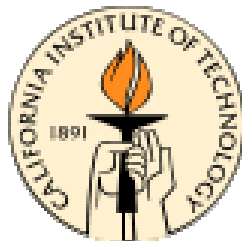
$$S_L = \sum_{m=1}^L H_2\left(\frac{1+v_m}{2}\right)$$



Computation of entanglement : finite XXZ spin chain

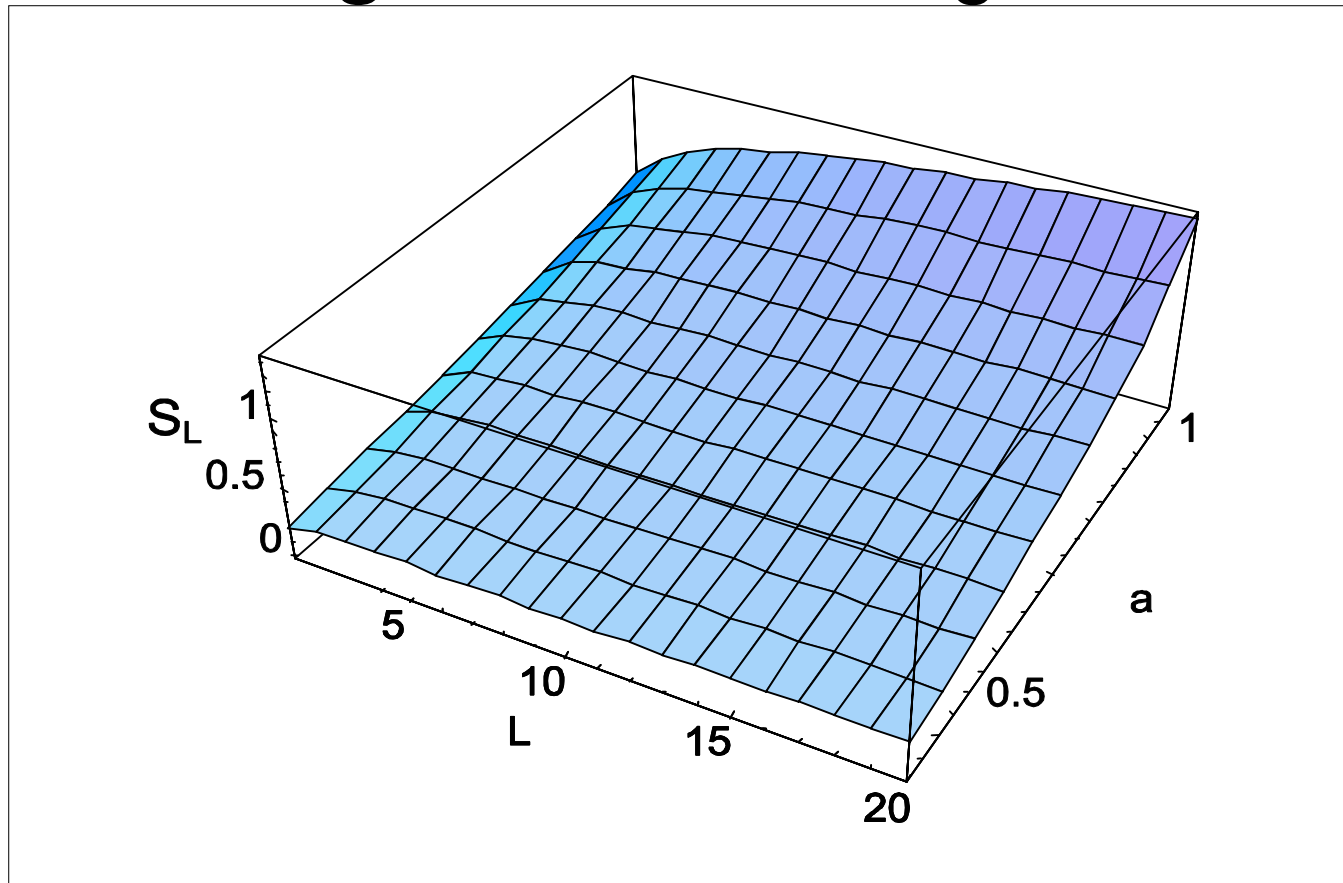
- Bethe ansatz for N=20 spins

$$|\Psi_g\rangle \rightarrow \rho_L \rightarrow S_L$$

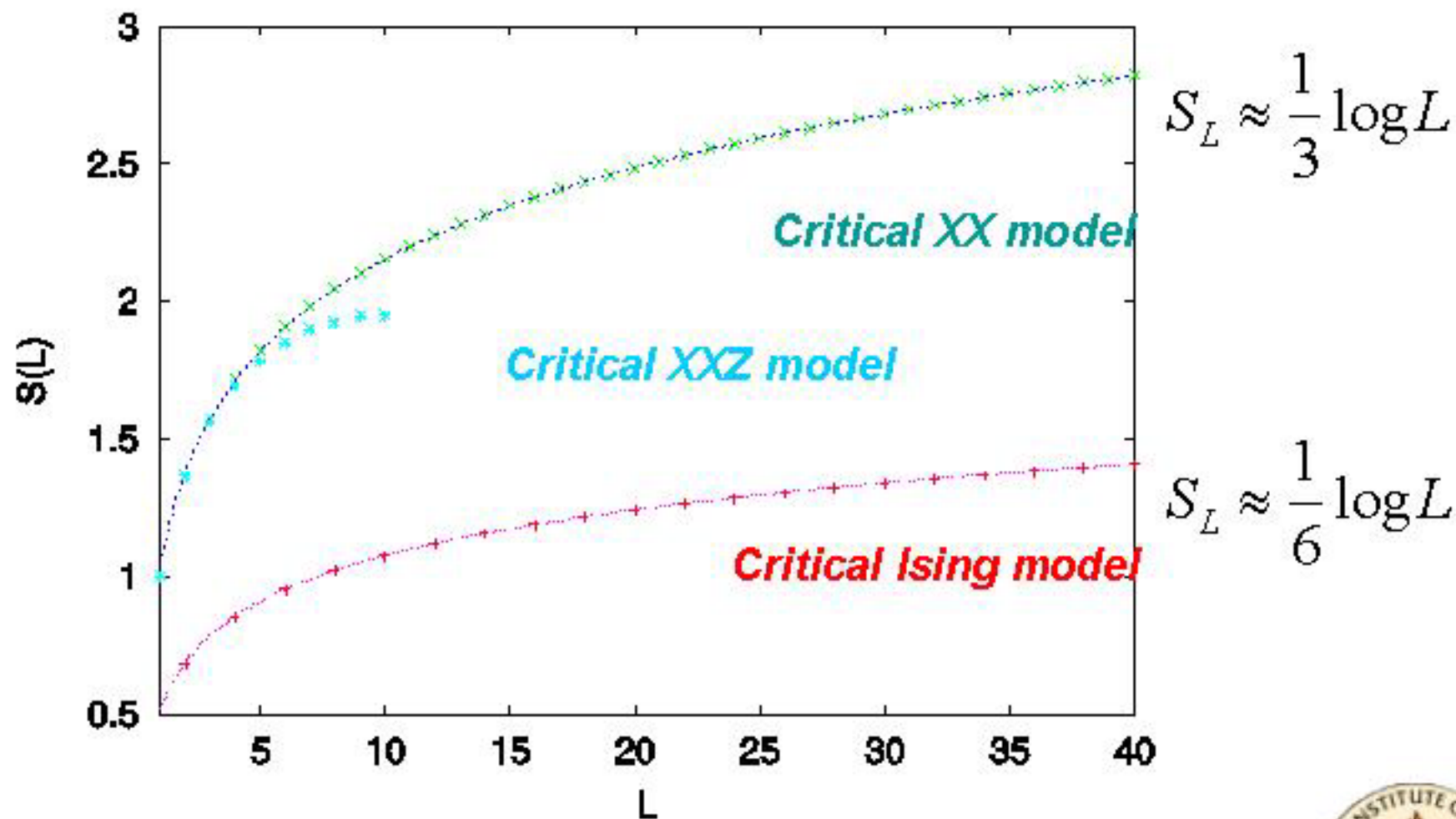


Non-critical entanglement

Ising chain with magnetic field



Critical entanglement



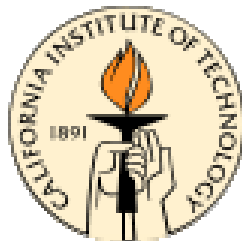
Non-critical versus critical ground-state entanglement

- *Non-critical* entanglement has a *saturation value*

$$S_L \leq S_{\max}$$

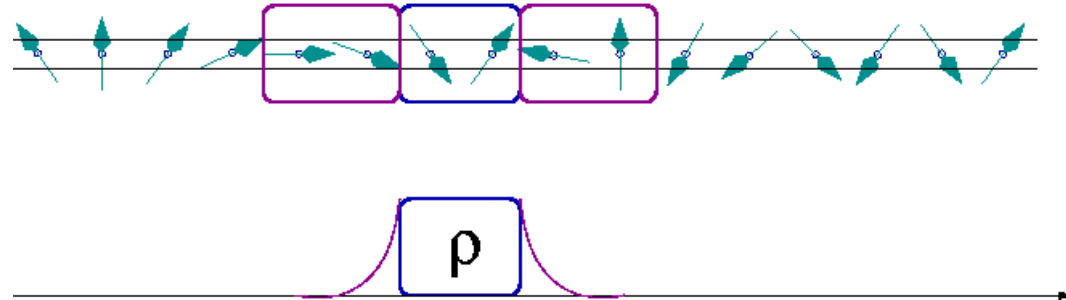
- *Critical* entanglement *diverges logarithmically* with the number L of spins

$$S_L \approx k \log L$$

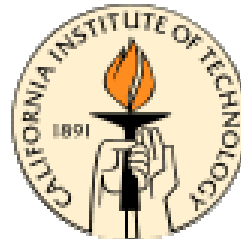
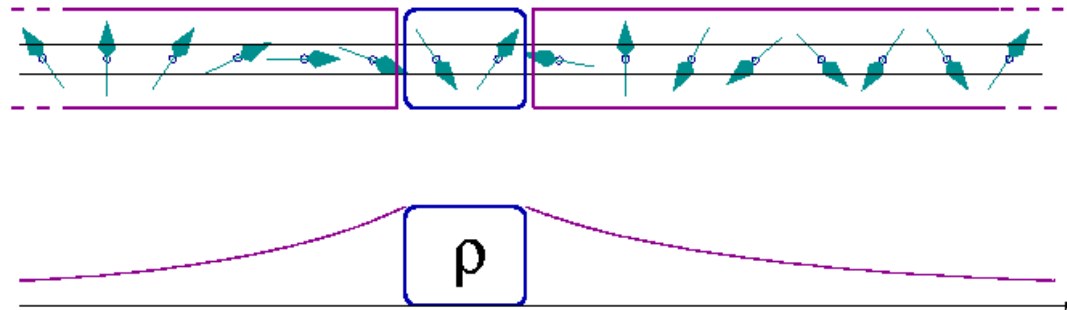


Density Matrix Renormalization Group Techniques

Non-critical entanglement is **semi-local**: The state of L qubits can be reconstructed by considering a few neighbors \rightarrow DMRG



Critical entanglement embraces the system at **all length scales**. It is not possible to construct ρ_L locally (its rank diverges) by DMRG



Connection to conformal field theory

- Geometric or fine-grained entropy

$$S_L \approx \frac{c + \bar{c}}{6} \log L$$

Holzhey, Larsen, Wilczek,
Nucl. Phys. B 424 (1994)

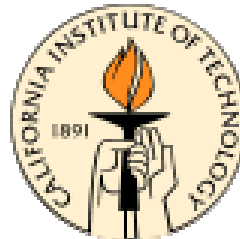
Srednicki, PRL 71 (1993)

Fiola, Preskill, Strominger,
Trivedi, PRD 50 (1994)

c is the **central charge** of the theory (holomorphic and antiholomorphic central charges)

$c_b = \bar{c}_b = 1$ for a **free boson** (XX model)

$c_f = \bar{c}_f = 1/2$ for a **free fermion** (Ising model)



Entanglement in 2D and 3D

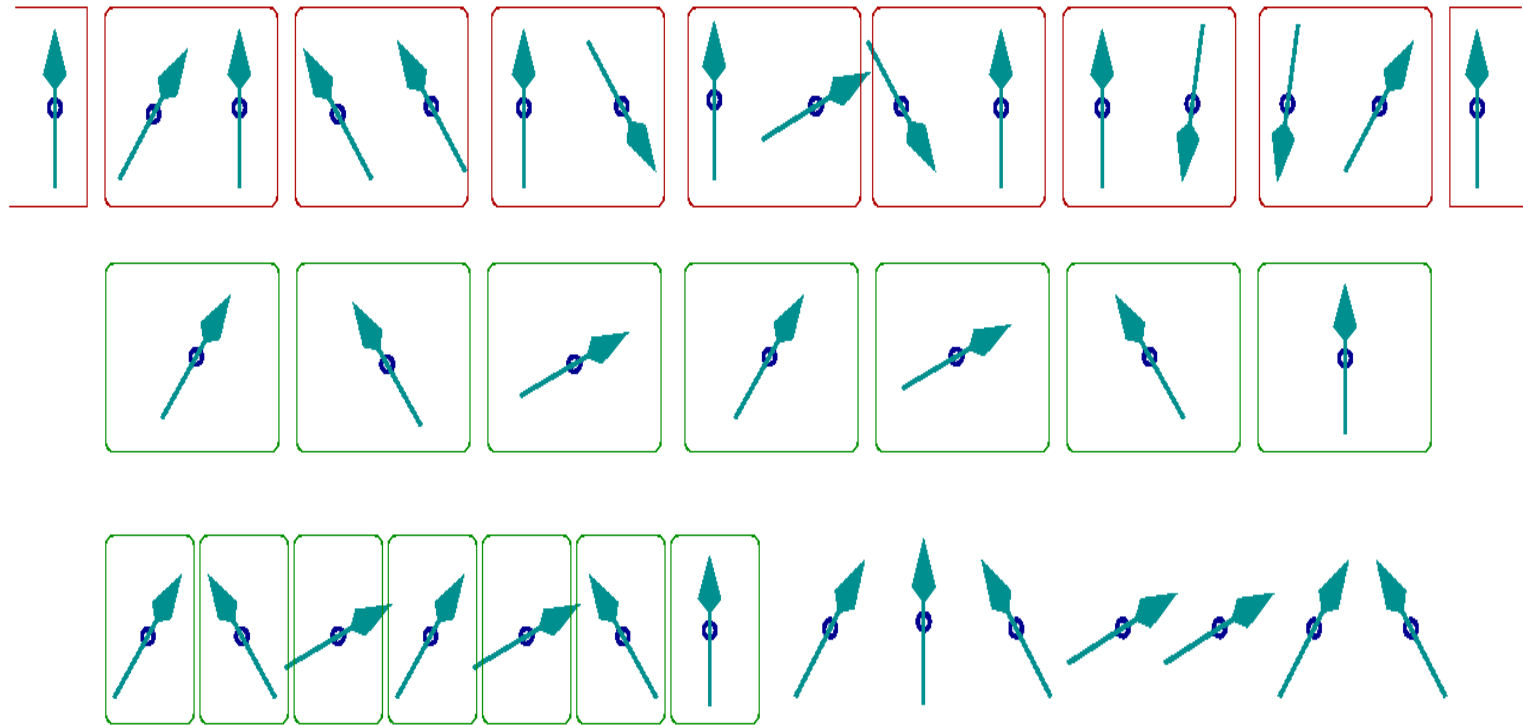
- We can export a result by Srednicki, PRL 71 (1993)

$$S_R \approx \kappa \Sigma(R)$$

The entanglement of a region R grows proportional to the size $\Sigma(R)$ of the boundary of R .



Renormalization Group flow



c-theorem

- The c-theorem establishes that the *central charge* can only *decrease* along the renormalization group flow

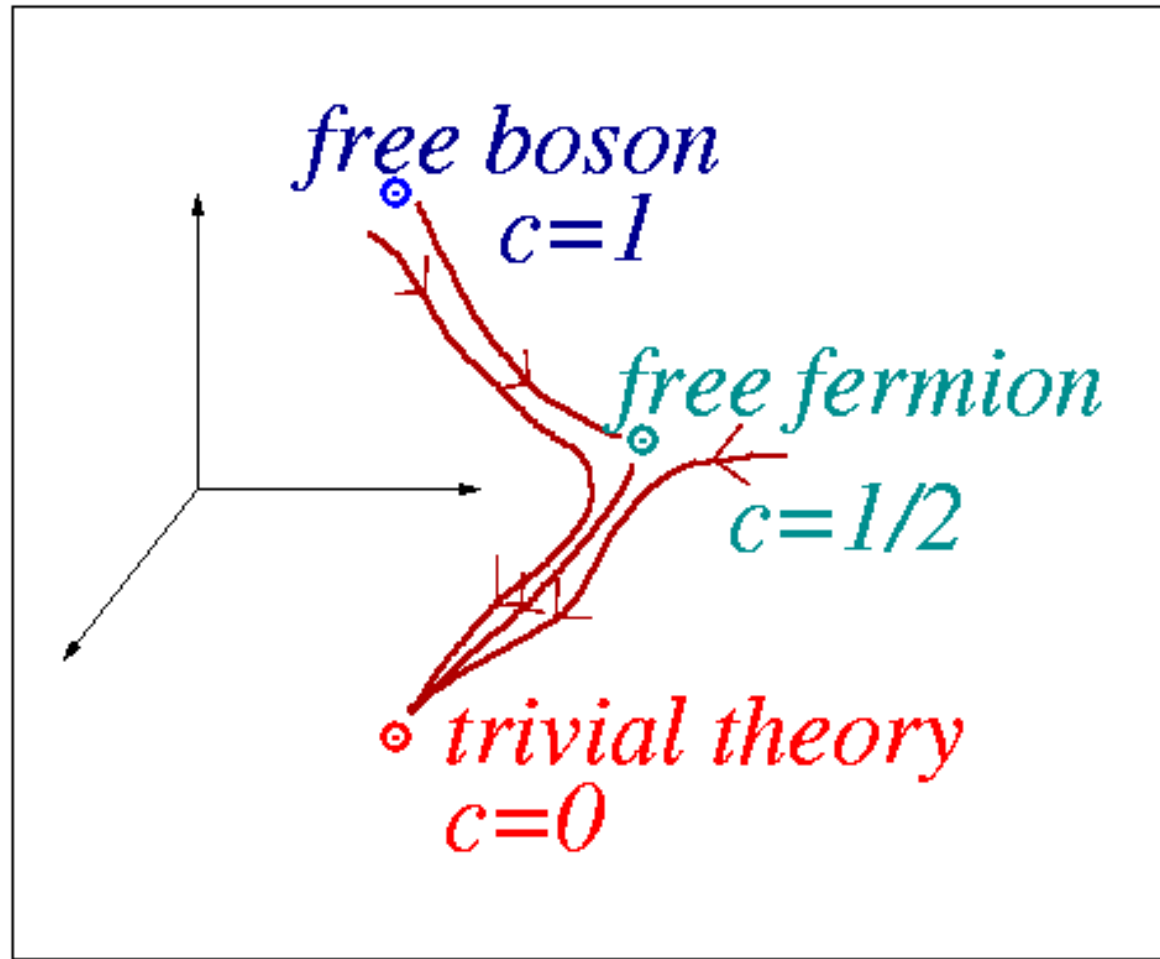
Zamolodchikov, JETP Lett 43 (1986)

Capelli, Friedan, Latorre, Nucl. Phys. B 352 (1991)

Forte, Latorre, Nucl. Phys. B 535 (1998)

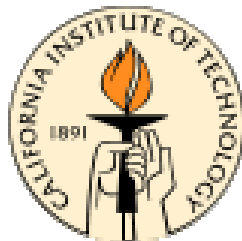


Renormalization Group flow and entanglement



Monotonicity of entanglement

- Under LOCC (local manipulation of a composite system)
- Along RG flow (change of scale)



Conclusions

Computation of entanglement in several
1D *critical* and *non-critical* spin models

- *Two distinctive forms of entanglement*

→ Breakdown of DMRG techniques

- *Connection to conformal field theory*

→ 2D and 3D entanglement

→ Monotonicity under RG transformations

