

# Optimal Transport in non-commutative probability

① Two views on heat flow:

\* Classical  $L^2$ -theory

- Gauss space  $L^2(\mathbb{R}^n, \gamma)$ .  $\gamma(dx) = \frac{1}{(2\pi)^{n/2}} e^{-\frac{|x|^2}{2}} dx$ .

- Dirichlet energy  $\mathcal{E}_r(\psi) = \int_{\mathbb{R}^n} |\nabla \psi(x)|^2 d\gamma(x)$ .

-> Gradient flow eq. for  $\mathcal{E}_r$  in  $L^2(\mathbb{R}^n, \gamma)$  is

the Ornstein-Uhlenbeck equation:  $\partial_t \psi = -\mathcal{L}\psi$ .

where  $-\mathcal{L}\psi(x) = \Delta \psi(x) - \langle x, \nabla \psi(x) \rangle$ .

\* optimal transport.

- 2-Wasserstein space  $(\mathcal{P}_2(\mathbb{R}^n), W_2)$

- Relative entropy  $\text{Ent}_\gamma(\mu) = \int f(x) \log p(x) d\mu(x)$ ,  $f = \frac{d\mu}{d\gamma}$

-> G.F. of  $\text{Ent}_\gamma$  w.r.t.  $W_2$  is the same O-U equation.

② Quantum Mechanical Perspective on  $L^2(\mathbb{R}^n, \gamma)$

$\bigoplus_{m \geq 0} (\mathbb{R}^n)^{\otimes_{\text{sym}} m}$  — particles.  $\cong \mathcal{F}_{\text{sym}}^n$ , Symmetric Fock space.

Wiener: Canonical Isometry:  $\mathcal{I}: L^2(\mathbb{R}^n, \gamma) \rightarrow \mathcal{F}_{\text{sym}}^n$ .

$\mathcal{I}: H_{m_1}(\langle \cdot, e_1 \rangle) \cdots H_{m_n}(\langle \cdot, e_n \rangle) \mapsto \mathcal{P}_{\text{sym}}(e_1^{\otimes m_1} \otimes \cdots \otimes e_n^{\otimes m_n})$

$\mathcal{N} = \mathcal{I} \circ \mathcal{L} \circ \mathcal{I}^{-1}$  is the boson number operator.

$\mathcal{N}h = mh$ .  $\forall h \in (\mathbb{R}^n)^{\otimes_{\text{sym}} m}$

### ③ Fermionic Counterpart

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• Anti-symmetric Fock space:  $\mathcal{F}_{\text{asym}}^n = \bigoplus_{m=0}^n (\mathbb{R}^n)^{\wedge m}$

• Fermionic <sup>Number</sup> Operator:  $Nh = mh, \quad \forall h \in (\mathbb{R}^n)^{\wedge m}$

Def. Let  $Q_1, \dots, Q_n$  be self-adjoint operators on a Hilbert space site.

$$\begin{cases} Q_i Q_j = -Q_j Q_i & \forall i \neq j \\ Q_i^2 = I & \forall i \end{cases}$$

Let  $C$  be the algebra  $\langle Q_1, \dots, Q_n \rangle$ . ~~called~~ called Clifford algebra.

Basis of  $C$  are  $Q_\alpha = Q_1^{\alpha_1} \dots Q_n^{\alpha_n}$  where  $\alpha_1, \dots, \alpha_n \in \{0, 1\}$ .

### ④ ~~W~~ Fermionic Wiener Isometry:

$$I: C \rightarrow \mathcal{F}_{\text{asym}}^n$$

$$Q_\alpha \mapsto e^{i\alpha_1} \wedge \dots \wedge e^{i\alpha_n}$$

### ④ - Analysis on the Clifford algebra

1. N.C. integration:  $\tau: C \rightarrow \mathbb{C}$ , by

$$\tau(Q_\alpha) = \begin{cases} 1 & \alpha = \underline{0} \\ 0 & \text{otherwise} \end{cases}$$

( $\tau$  is a positive linear functional satisfying  $\tau(AB) = \tau(BA)$ .)

Remark.  $\tau$  induces a scalar product

$$\langle A, B \rangle_{L^2} = \tau(A^* B)$$

2. N.C. Differentiation: define  $\partial_j: \mathcal{C} \rightarrow \mathcal{C}$  by

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$$\partial_j A = \frac{1}{i} (Q_j A - \Gamma(A) Q_j), \quad \text{where } \Gamma(Q_\alpha) = (-1)^{|\alpha|} Q_\alpha.$$

3. Def: (Clifford Dirichlet Form, Gauss)

$$\mathcal{E}: \mathcal{C} \times \mathcal{C} \rightarrow \mathcal{C}$$

$$\mathcal{E}(A, B) = \sum_{j=1}^n \langle \partial_j A, \partial_j B \rangle_{\mathcal{L}^2}$$

4. Prop. The associated operator is the unbiased operator.  
i.e.  $\mathcal{E}(A, B) = \langle \mathcal{L}A, B \rangle_{\mathcal{L}^2}$ , where  $\mathcal{L} = I^{\dagger} \circ \mathcal{N} \circ I$

(5) General setup.

•  $(A, \tau)$  non-commutative prob space, i.e.

•  $A$  is a  $\ast$ -algebra of operators on a Hilbert space (finite dim)

•  $\tau: A \rightarrow \mathbb{C}$  trace (positive linear functional s.t.  $\tau(AB) = \tau(BA)$ )

•  $\partial_j: A \rightarrow A_j$ ,  $j=1, \dots, n$ , operator of the form

$$\partial_j A = Q_j \Gamma_j(A) - \Gamma_j(A) Q_j, \quad \text{where}$$

$(A_j, \tau_j)$  non-comm. prob. space,  $Q_j \in A_j$ ,  $\Gamma_j, \tau_j: A \rightarrow A_j$   
algebra hom.

# ⑥. Examples

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1). Clifford algebra

2). Continuous time Markov chains with transition rates  $Q(x,y)$  on reversible inv. measure  $\pi$ .

$$E(\psi) = \frac{1}{2} \sum_{x,y} (\psi(x) - \psi(y))^2 Q(x,y) \pi(x)$$

$$L\psi(x) = \sum_y Q(x,y) (\psi(y) - \psi(x))$$

3). Lindblad equations (open quantum systems)

⑦. Recall: Bonami - Brierley formula for  $\rho_0, \rho_1 \in \mathcal{P}(\mathbb{R}^n)$

$$W_2^2(\rho_0, \rho_1) = \inf \left\{ \int_0^1 \int_{\mathbb{R}^n} |\nabla \psi_t(x)|^2 d\rho_t(x) dt \right\}$$

$$\text{s.t. } \partial_t \rho_t + \nabla \cdot (\rho_t \nabla \psi_t) = 0, \quad \left. \begin{array}{l} \rho_{t=0} = \rho_0 \\ \rho_{t=1} = \rho_1 \end{array} \right\}$$

Problem: How to define " $\rho * \partial_j \psi$ "?

In  $\mathbb{R}^n$ , heat equation as continuity equation.

$$\partial_t \rho = \Delta \rho \iff \begin{cases} \partial_t \rho + \nabla \cdot (\rho \nabla \psi) = 0 \\ \psi = -\log \rho \end{cases}$$

use chain rule:  $\nabla \rho = \rho \cdot \nabla \log \rho$

# ⑧. Non-Commutative Chain Rule

We have:  $\partial_j(AB) = \partial_j(A) \cdot B + A \cdot \partial_j(B)$  (Assume  $\ell = r = \text{id}$ ).

Induction:  $\partial_j(A^n) = \sum_{k=0}^{n-1} A^k \cdot \partial_j A \cdot A^{n-k-1}$

Set  $\rho = A^n$ :  $\partial_j \rho = \sum_{k=0}^{n-1} \rho^{\frac{k}{n}} \partial_j \rho^{\frac{1}{n}} \rho^{\frac{n-k-1}{n}}$

Then  $\partial_j p = \frac{1}{n} \sum_{k=0}^{n-1} p^{\frac{k}{n}} \frac{\partial_j (p^{\frac{k}{n}} - I)}{\frac{1}{n}} p^{1 - \frac{k+1}{n}}$

(5)

Let  $n \rightarrow \infty$ ,

Then  $\partial_j p = \int_0^1 p^x \cdot \partial_j \log p \cdot p^{1-x} dx$

9 Def. Set  $\mathcal{P}(A) = \{p \in A : p \geq 0, \text{Tr}(p) = 1\}$  (prob. densities)

Def ("2-Wasserstein metric")

For  $p_0, p_1 \in \mathcal{P}(A)$ , set

$$W^2(p_0, p_1) = \inf \left\{ \int_0^1 \langle p_t \times \partial_t \psi_t, \partial_t \psi_t \rangle_{L^2} dt \right\}$$

$$\left. \begin{aligned} & \partial_t p + \sum \partial_j^* (p_t \times \partial_j \psi_t) = 0 \\ & p_{t=0} = p_0, p_{t=1} = p_1 \end{aligned} \right\}$$

where  $p \times A = \int_0^1 \log(p)^{1-x} A r_j(p)^x dx$ .

10 Prop.  $W$  is associated to Riemannian structure on  $\mathcal{P}(A)$ .

Thm. Set  $\text{Ent}(p) = \int \tau(p \log p)$  (von Neumann entropy)

The grad. flow equation for  $\text{Ent}$  w.r.t  $W$  is

$$\partial_t p = -\mathcal{L} p, \text{ where } \mathcal{L} = \sum_{j=1}^n \partial_j^* \partial_j.$$

In eq. 1). (Carlen - M.)

eq. 2). (Mielke, Chow - Hongy - Li - Zhou, M.)

eq. 3). Mielke.

(11). Then (Carver - Marcus)

(6)

If  $(A, \tau, (\partial_j))$ , satisfies a log-Sobolev. iner.

$$\text{i.e. } \text{Ent}(P) \leq \frac{1}{2K} \sum_j \langle \partial_j P, \partial_j \log P \rangle$$

Then non-comm. Talagrand ~~is~~ ineq. holds.

$$\text{(i.e. } W(p, I)^2 \leq \frac{2}{K} \text{Ent}(p) \quad \forall p \text{)}$$

(analogue of Otto-Villani).